

STATUS of $\bar{K}N$ and $\bar{K}NN$ INTERACTIONS

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KEYWORDS:

- **LOW-ENERGY QCD** with **STRANGE QUARKS** realized as an **EFFECTIVE FIELD THEORY**: **SU(3)** octet of pseudoscalar Nambu-Goldstone bosons coupled to the baryon octet
- **Update on $\bar{K}N$ and $\bar{K}NN$ interactions**
Scattering lengths, quasibound states, two-poles scenario, ...



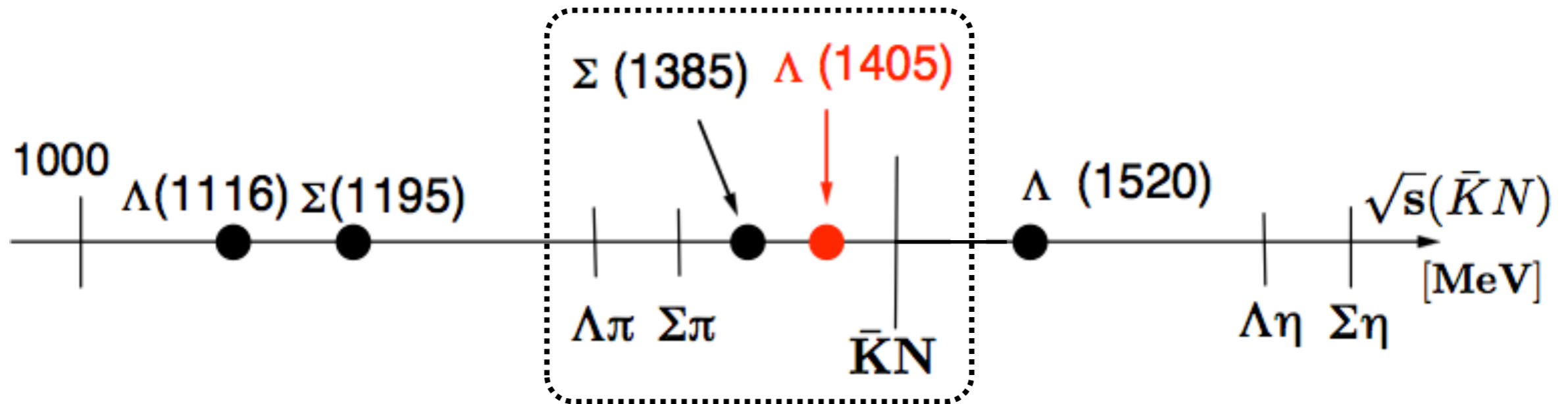
BASIC ISSUES

- **Strange quarks** are intermediate between “**light**” and “**heavy**”:
 - ▶ interplay between **spontaneous** and **explicit chiral symmetry breaking** in low-energy QCD
- Testing ground: high-precision **antikaon-nucleon** threshold physics
 - ▶ strongly **attractive** low-energy $\bar{K}N$ interaction
- Nature and structure of $\Lambda(1405)$ ($B = 1$, $S = -1$, $J^P = 1/2^-$)
 - ▶ **three-quark** valence structure vs. “**molecular**” meson-baryon system ?
- Quest for quasi-bound **antikaon-nuclear** systems ?
- Role of **strangeness** in dense baryonic matter ?
 - ▶ new constraints from **neutron stars**



LOW-ENERGY $\bar{K}N - \pi Y$ SYSTEMS

- **Poles and thresholds:**



$\Lambda(1405)$ resonance 27 MeV below threshold:

▶ chiral perturbation theory **NOT** applicable

- Strategy:

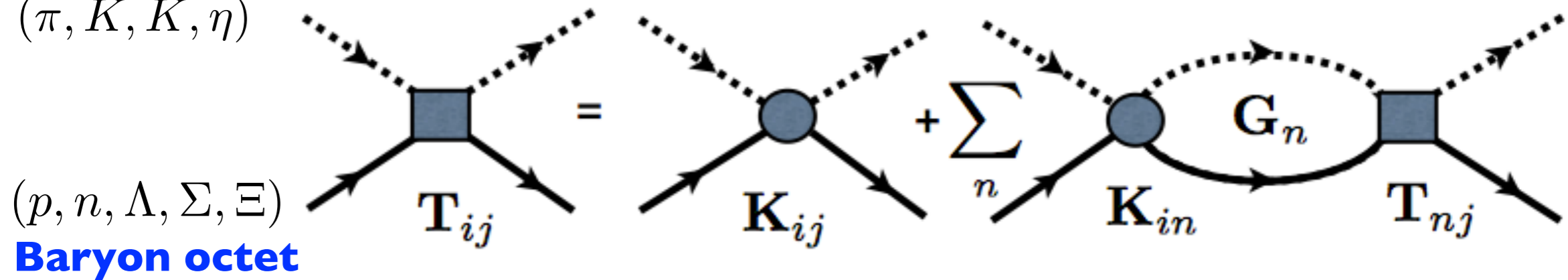
Non-perturbative Coupled-Channels Dynamics
based on **Chiral SU(3) Effective Lagrangian**



CHIRAL SU(3) DYNAMICS with **COUPLED CHANNELS**

Pseudoscalar meson octet

(π, K, \bar{K}, η)



$$\mathbf{T} = \mathbf{K} + \mathbf{K} \mathbf{G} \mathbf{T} = (\mathbf{1} - \mathbf{K} \mathbf{G})^{-1} \mathbf{K}$$

$$\mathbf{T}_{ij}(p', p, \sqrt{s}) = \mathbf{K}_{ij}(p', p, \sqrt{s}) + \sum_n \int \frac{d^4 q}{(2\pi)^4} \mathbf{K}_{in}(p', q, \sqrt{s}) \mathbf{G}_n(q, \sqrt{s}) \mathbf{T}_{nj}(q, p, \sqrt{s})$$

Kernel \mathbf{K}_{ij} from

CHIRAL SU(3) EFFECTIVE MESON-BARYON LAGRANGIAN



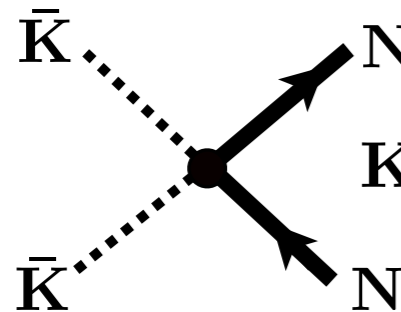
CHIRAL SU(3) COUPLED CHANNELS DYNAMICS

$$T_{ij} = K_{ij} + \sum_n K_{in} G_n T_{nj}$$

- Leading s-wave $I = 0$ meson-baryon interactions (Tomozawa-Weinberg)

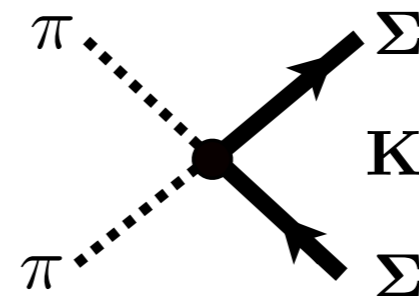
Note: **ENERGY DEPENDENCE** characteristic of Nambu-Goldstone Bosons

$$|1\rangle = |\bar{K}N, I = 0\rangle$$



$$K_{11} = \frac{3}{2f_K^2} (\sqrt{s} - M_N)$$

$$|2\rangle = |\pi\Sigma, I = 0\rangle$$



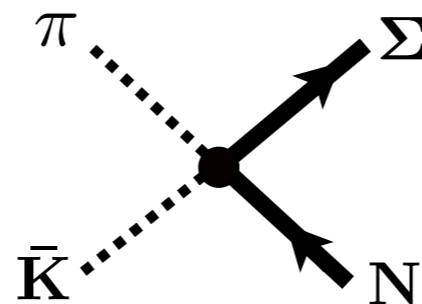
$$K_{22} = \frac{2}{f_\pi^2} (\sqrt{s} - M_\Sigma)$$

- driving interactions individually **strong** enough to produce

▶ $\bar{K}N$ **bound state**

▶ $\pi\Sigma$ **resonance**

- strong** channel coupling
 $12 \leftrightarrow 21$:



$$K_{12} = \frac{-1}{2f_\pi f_K} \sqrt{\frac{3}{2}} \left(\sqrt{s} - \frac{M_\Sigma + M_N}{2} \right)$$

$$f_\pi = 92.4 \pm 0.3 \text{ MeV}$$

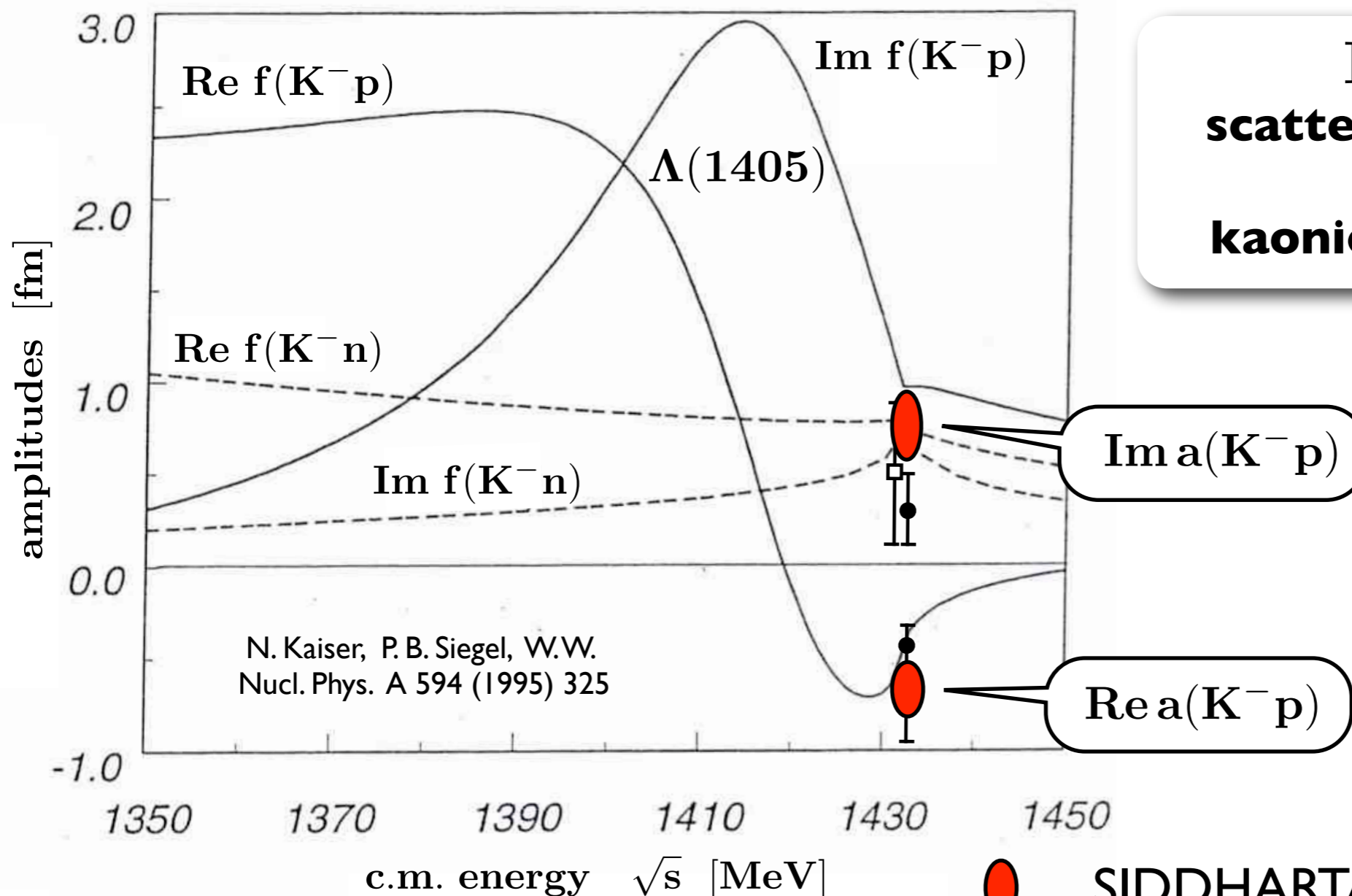
$$f_K = 110.0 \pm 0.9 \text{ MeV}$$



$\bar{K} N$ AMPLITUDES - past and present -

CHIRAL SU(3) EFFECTIVE FIELD THEORY with COUPLED CHANNELS

leading order (Tomozawa - Weinberg) terms



$K^- p$
scattering length
from
kaonic hydrogen

KEK

M. Iwasaki et al.
Phys. Rev. Lett.
78 (1997) 3067

DEAR @ LNF

G. Beer et al.
Phys. Rev. Lett.
94 (2005) 212302



SIDDHARTA @ LNF

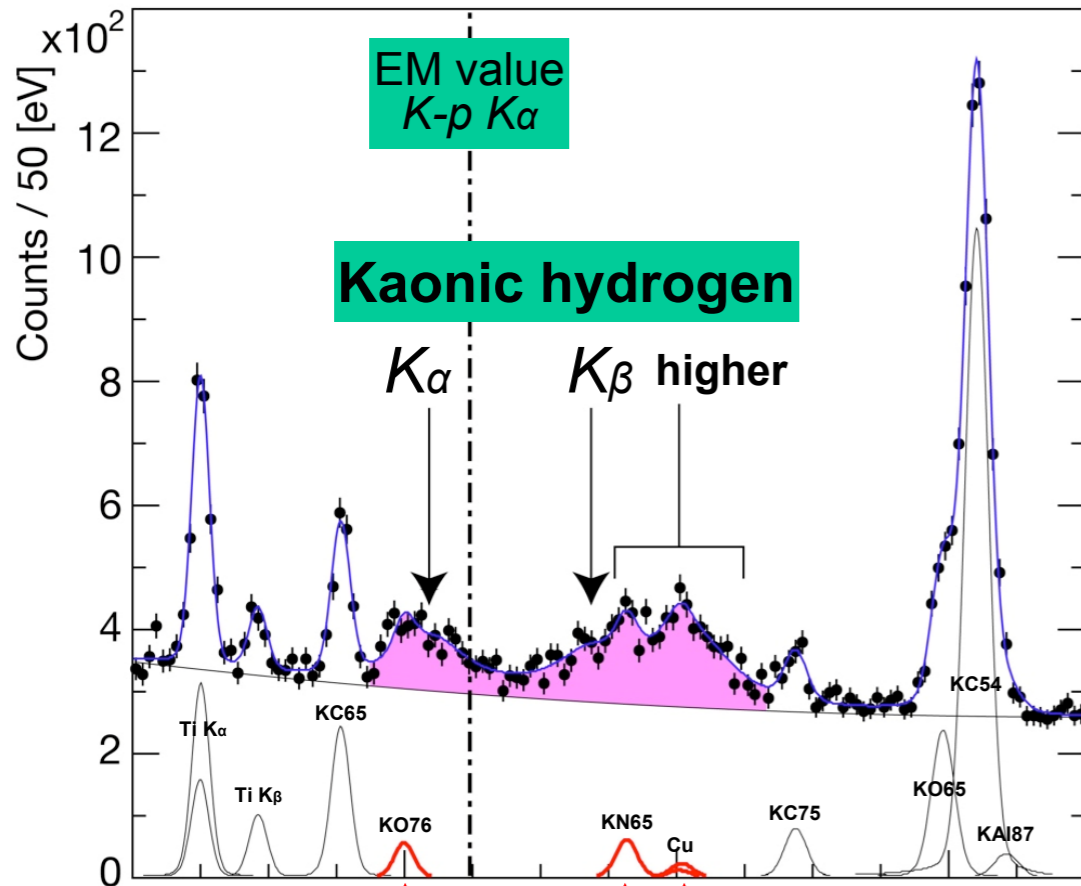
M. Bazzi et al. : Phys. Lett. B 704 (2011) 113



NEWS from SIDDHARTA

● **Kaonic hydrogen** precision data

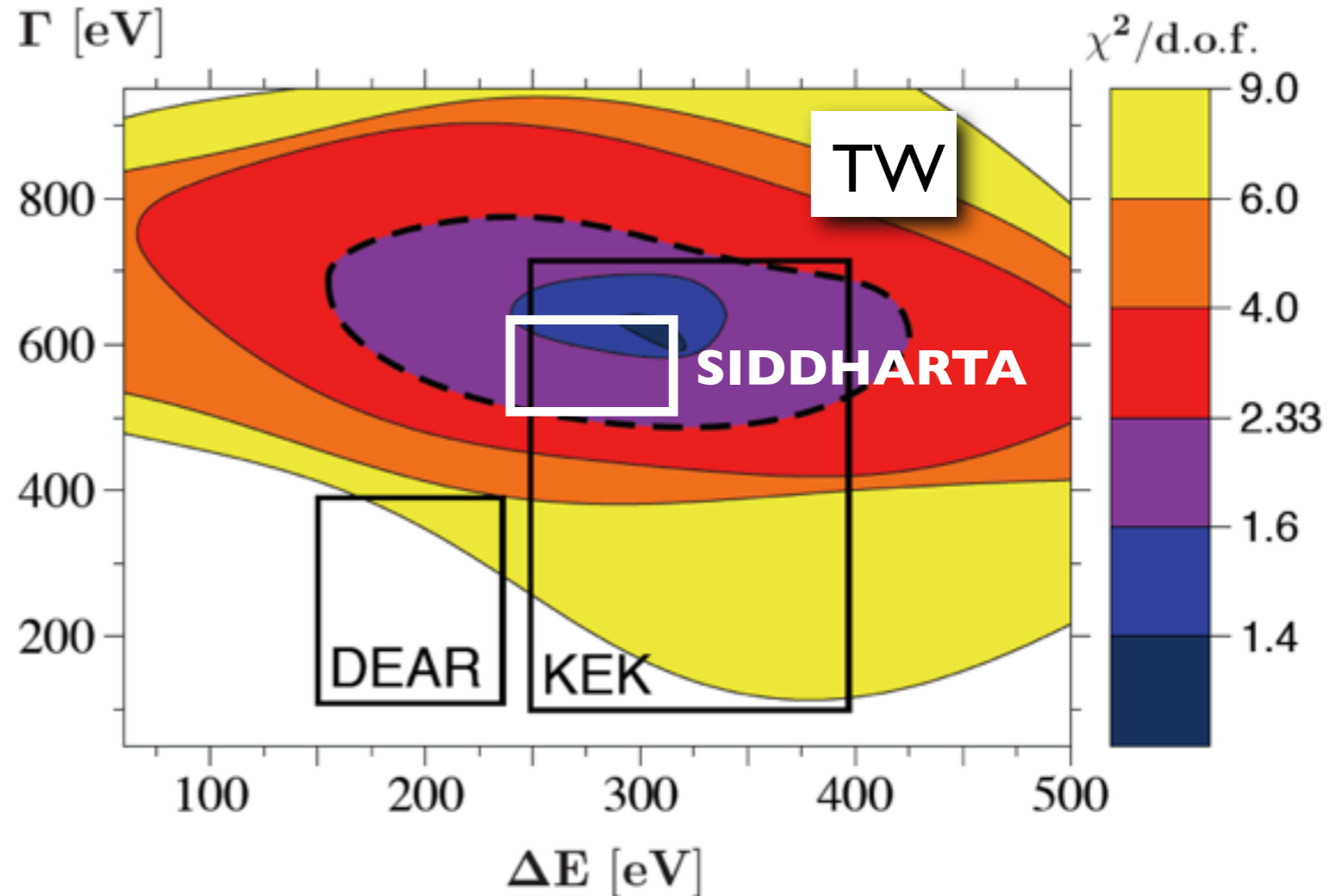
● strong interaction shift and width:



M. Bazzi et al.
Phys. Lett. B 704 (2011) 113

$$\Delta E = 283 \pm 36 (stat) \pm 6 (syst) \text{ eV}$$

$$\Gamma = 541 \pm 89 (stat) \pm 22 (syst) \text{ eV}$$



R. Nißler PhD thesis (2008)

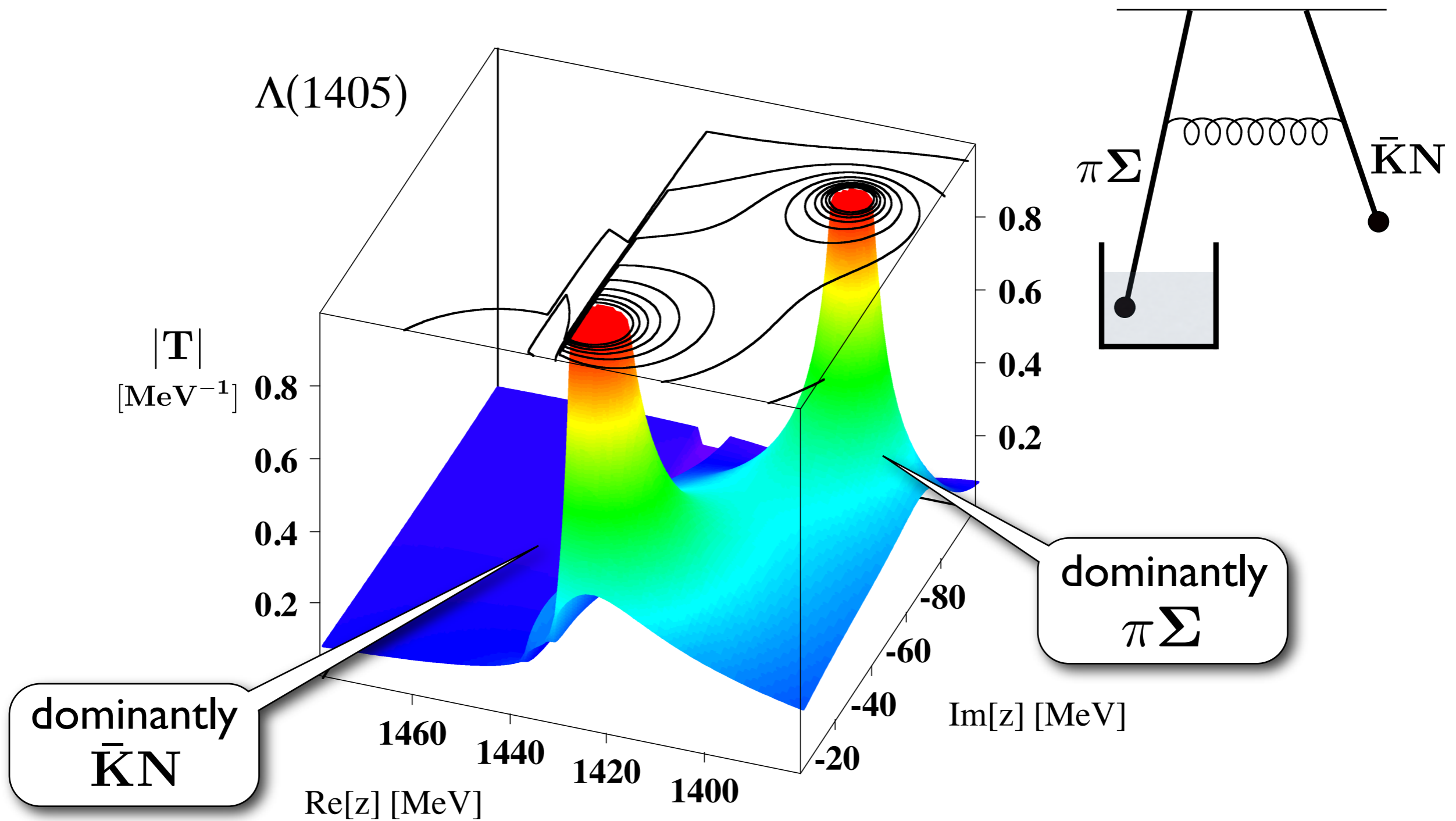
● **Theory:**
leading order
(Tomozawa - Weinberg)

B. Borasoy, R. Nißler, W.W. Eur. Phys. J. A25 (2005) 79

B. Borasoy, U.-G. Meißner, R. Nißler PRC74 (2006) 055201



The TWO POLES scenario



D. Jido et al., Nucl. Phys. A723 (2003) 205

T. Hyodo, W.W.: Phys. Rev. C 77 (2008) 03524

T. Hyodo, D. Jido : Prog. Part. Nucl. Phys. 67 (2012) 55



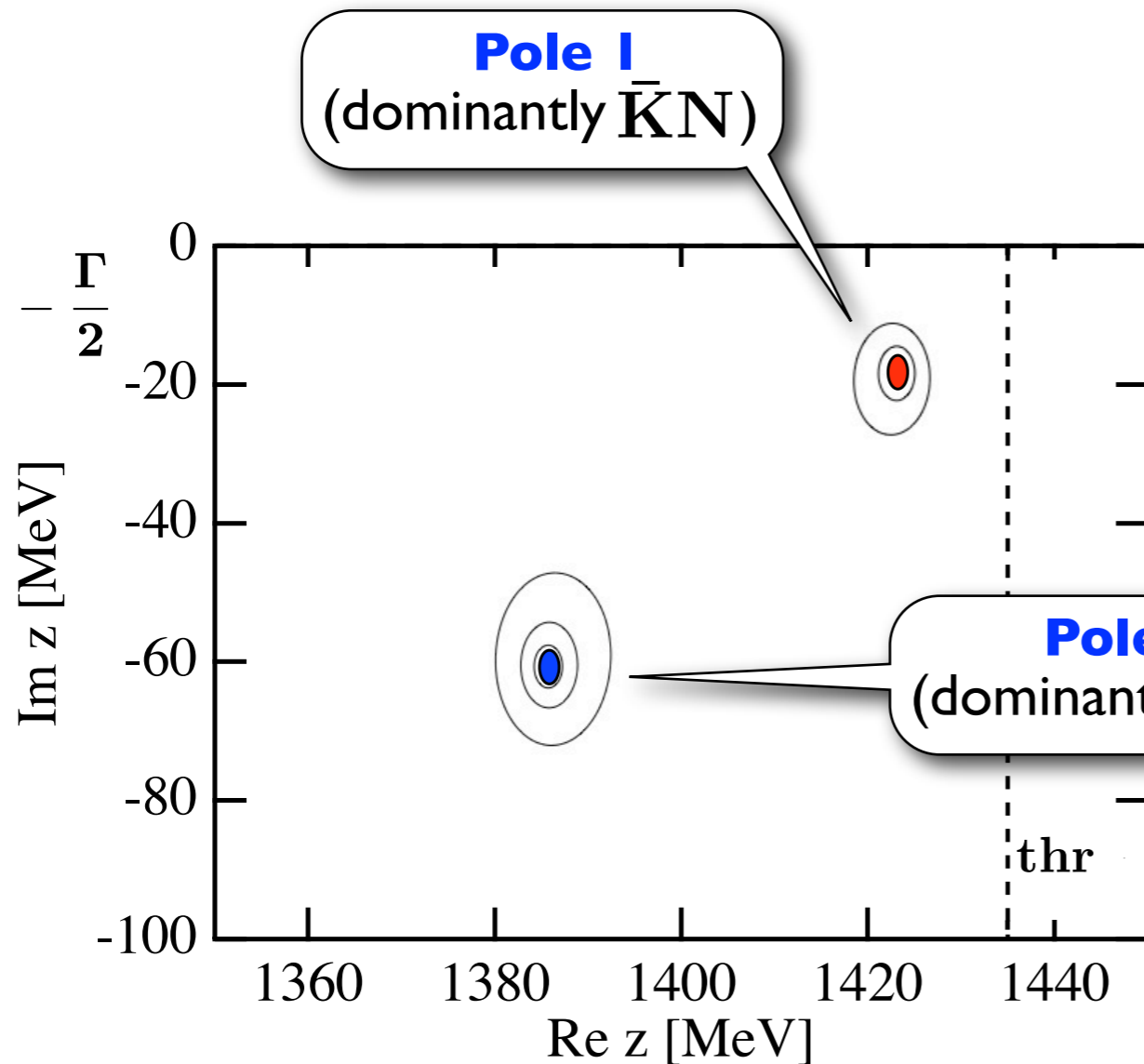
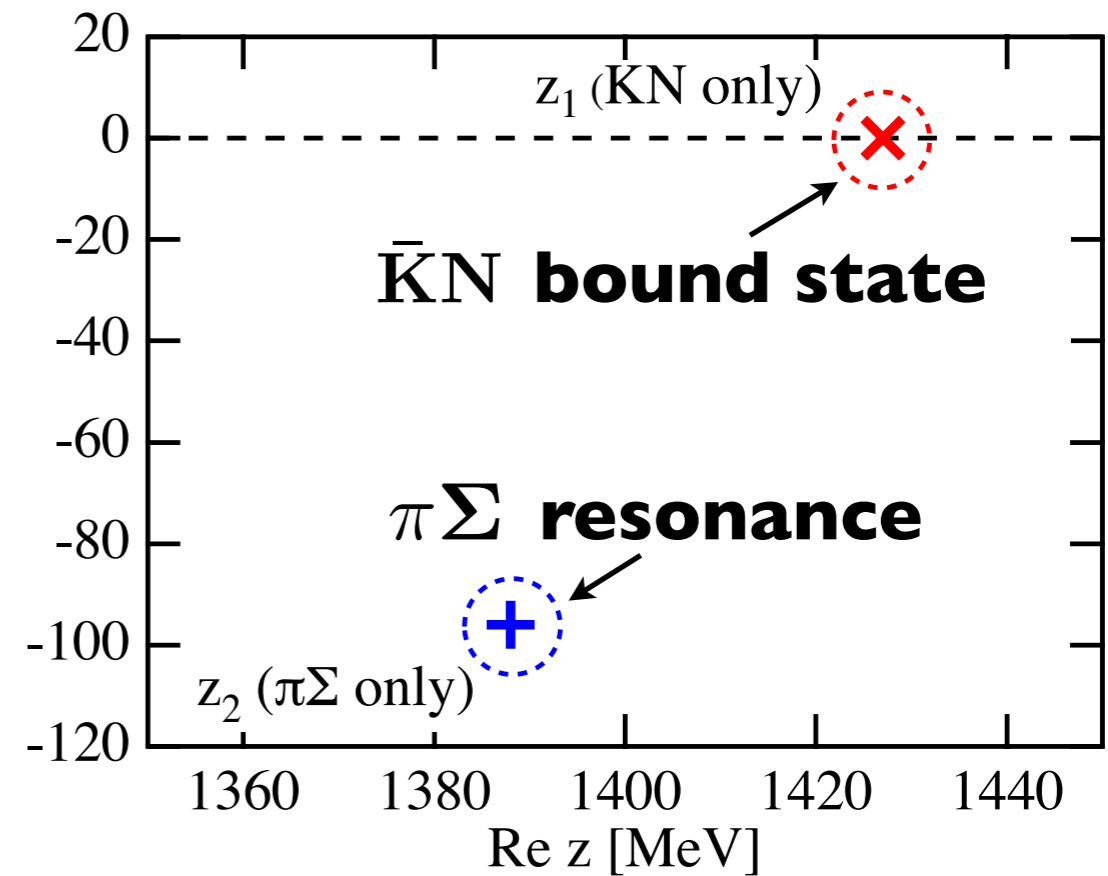
The TWO POLES scenario

T. Hyodo, W.W., Phys. Rev. C77 (2008) 03524

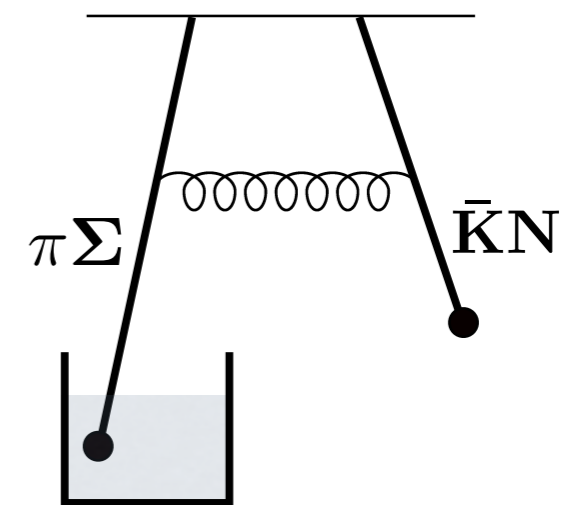
- Singularities of $\bar{K}N$ amplitude in the complex energy plane

starting point:

no channel coupling

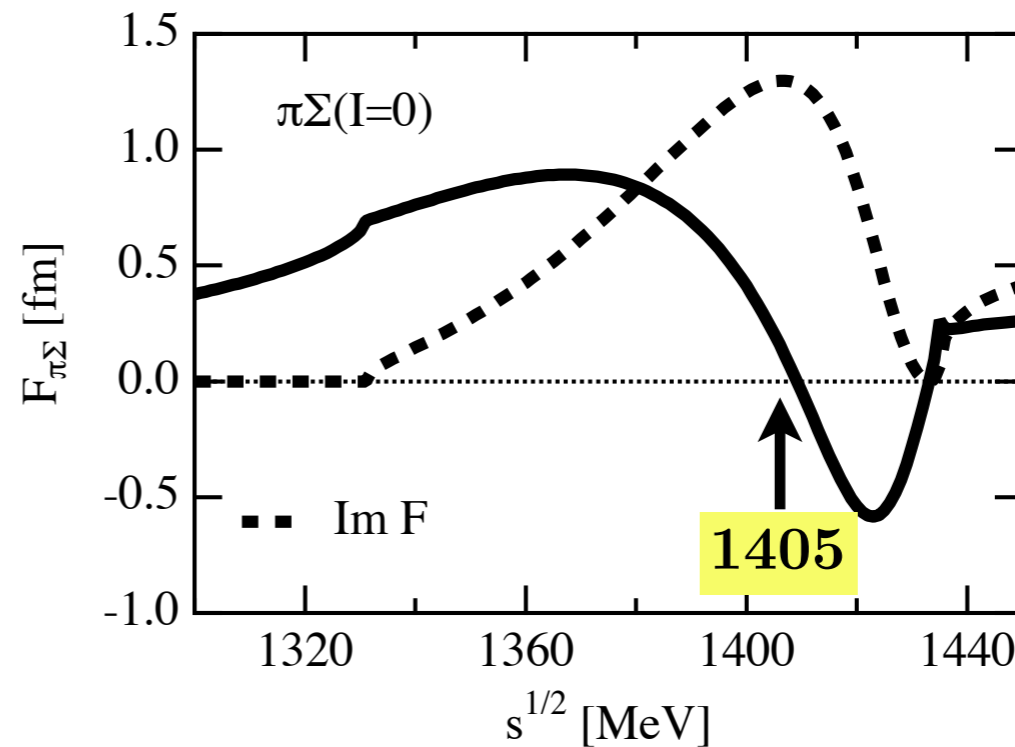
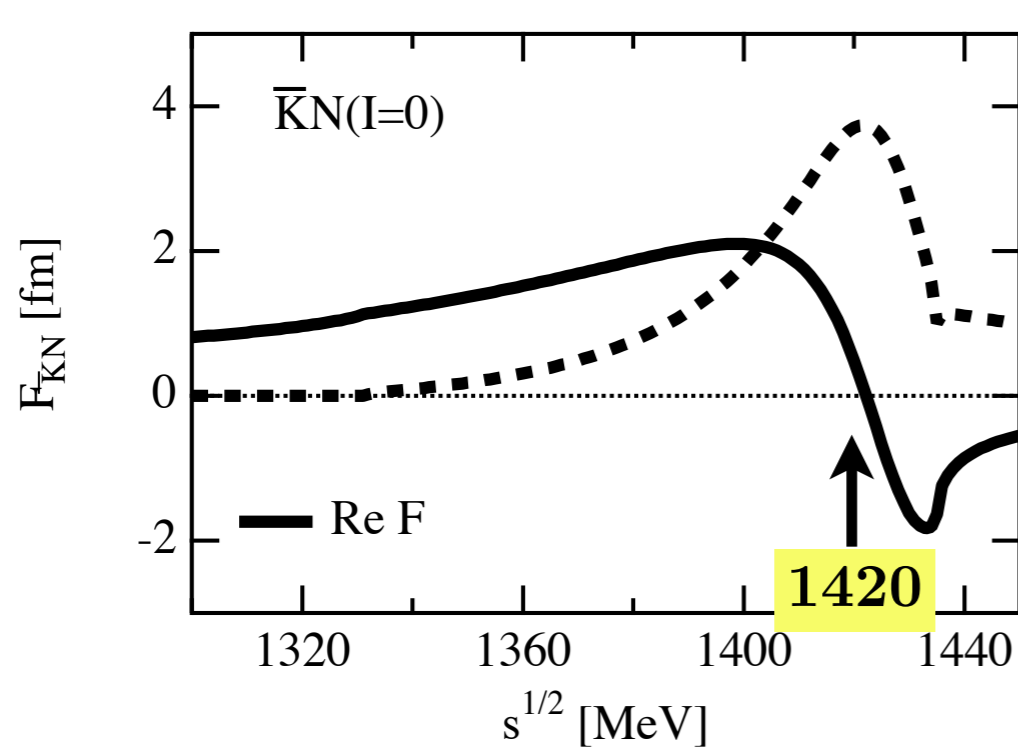


channel coupling at work



The TWO POLES scenario (contd.)

- $\bar{K}N$ and $\pi\Sigma$ amplitudes



T. Hyodo, W.W.: Phys. Rev. C77 (2008) 03524

- ▶ **Note difference** in spectral maxima of $\bar{K}N$ and $\pi\Sigma$

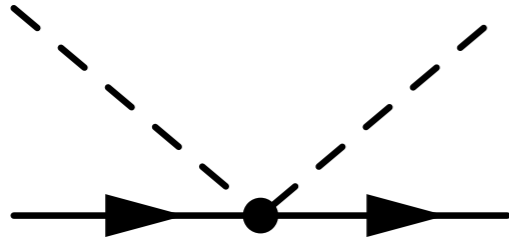
D. Jido et al., NP A725 (2003) 263

- ▶ Equivalent $\bar{K}N$ **effective interaction** should produce quasibound state at **1420 MeV (not 1405 MeV)**



CHIRAL SU(3) COUPLED CHANNELS DYNAMICS:

- NLO hierarchy of driving terms -

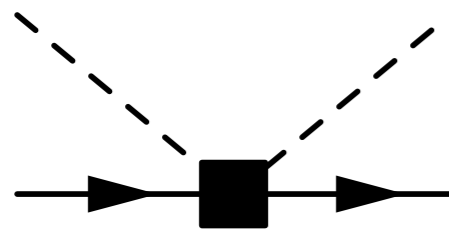
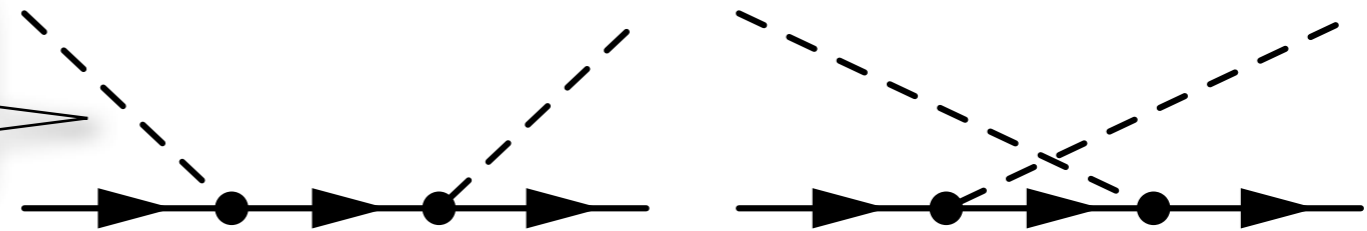


leading order (**W**einberg-**T**omozawa) terms
input: physical pion and kaon decay constants

direct and crossed **B**orn terms
input: axial vector constants
 D and F from hyperon beta decays

$$g_A = D + F = 1.26$$

$$\mathcal{L}_1^{MB} = \text{Tr} \left(\frac{D}{2} (\bar{B} \gamma^\mu \gamma_5 \{u_\mu, B\}) + \frac{F}{2} (\bar{B} \gamma^\mu \gamma_5 [u_\mu, B]) \right)$$



next-to-leading order (**NLO**)
input: 7 s-wave low-energy constants $\mathcal{O}(p^2)$

$$\begin{aligned} \mathcal{L}_2^{MB} = & b_D \text{Tr}(\bar{B} \{\chi_+, B\}) + b_F \text{Tr}(\bar{B} [\chi_+, B]) + b_0 \text{Tr}(\bar{B} B) \text{Tr}(\chi_+) \\ & + d_1 \text{Tr}(\bar{B} \{u^\mu, [u_\mu, B]\}) + d_2 \text{Tr}(\bar{B} [u^\mu, [u_\mu, B]]) \\ & + d_3 \text{Tr}(\bar{B} u_\mu) \text{Tr}(u^\mu B) + d_4 \text{Tr}(\bar{B} B) \text{Tr}(u^\mu u_\mu), \end{aligned}$$



CHIRAL SU(3) COUPLED CHANNELS DYNAMICS

(contd.)

The diagram illustrates the Dyson equation for the transition amplitude T_{ij} . On the left, a square box labeled T_{ij} has two solid lines entering from the bottom and two dashed lines exiting from the top. This is equal to a sum of two terms. The first term is a circle labeled K_{ij} with the same external lines. The second term is a sum over n of a diagram where a circle labeled K_{in} is connected to a square box labeled T_{nj} by a dashed loop labeled G_n . The external lines are the same as in the first term.

$$T_{ij} = K_{ij} + \sum_n K_{in} G_n T_{nj}$$

- **channels:** $K^-p, \bar{K}^0n, \pi^0\Sigma^0, \pi^+\Sigma^-, \pi^-\Sigma^+, \pi^0\Lambda, \eta\Lambda, \eta\Sigma^0, K^+\Xi^-, K^-\Xi^0$
- **loop integrals** (with meson-baryon Green functions) using dimensional regularization:

$$\tilde{G}(q^2) = \int \frac{d^d p}{(2\pi)^d} \frac{i}{[(q-p)^2 - M_B^2 + i\epsilon][p^2 - m_\phi^2 + i\epsilon]}$$

- finite parts including **subtraction constants** $a(\mu)$:

$$G(q^2) = a(\mu) + \frac{1}{32\pi^2 q^2} \left\{ q^2 \left[\ln\left(\frac{m_\phi^2}{\mu^2}\right) + \ln\left(\frac{M_B^2}{\mu^2}\right) - 2 \right] + (m_\phi^2 - M_B^2) \ln\left(\frac{m_\phi^2}{M_B^2}\right) - 8\sqrt{q^2} |\mathbf{q}_{cm}| \operatorname{artanh}\left(\frac{2\sqrt{q^2} |\mathbf{q}_{cm}|}{(m_\phi + M_B)^2 - q^2}\right) \right\}$$



UPDATED ANALYSIS of K^-p THRESHOLD PHYSICS

Y. Ikeda, T. Hyodo, W.W. Physics Letters B 706 (2011) 63 Nucl. Phys. A 881 (2012) 98

- Chiral SU(3) coupled-channels dynamics
Tomozawa-Weinberg + **Born** terms + **NLO**

kaonic hydrogen shift & width	theory (NLO)	exp.
ΔE (eV)	306	$283 \pm 36 \pm 6$
Γ (eV)	591	$541 \pm 89 \pm 22$
threshold branching ratios		(SIDDHARTA)
$\frac{\Gamma(K^-p \rightarrow \pi^+\Sigma^-)}{\Gamma(K^-p \rightarrow \pi^-\Sigma^+)}$	2.37	2.36 ± 0.04
$\frac{\Gamma(K^-p \rightarrow \pi^+\Sigma^-, \pi^-\Sigma^+)}{\Gamma(K^-p \rightarrow \text{all inelastic channels})}$	0.66	0.66 ± 0.01
$\frac{\Gamma(K^-p \rightarrow \pi^0\Lambda)}{\Gamma(K^-p \rightarrow \text{neutral states})}$	0.19	0.19 ± 0.02
scattering length (fm)	$\text{Re } a(K^-p) = -0.65 \pm 0.10$	$\text{Im } a(K^-p) = 0.81 \pm 0.15$

best fit achieved with $\chi^2/d.o.f. \simeq 0.9$



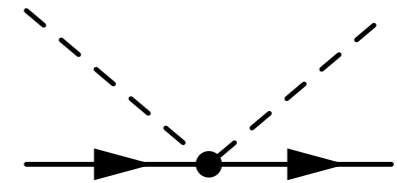
UPDATED ANALYSIS of K^-p THRESHOLD PHYSICS with **SIDDHARTA** constraints

Y. Ikeda, T. Hyodo, W.W. Physics Letters B 706 (2011) 63

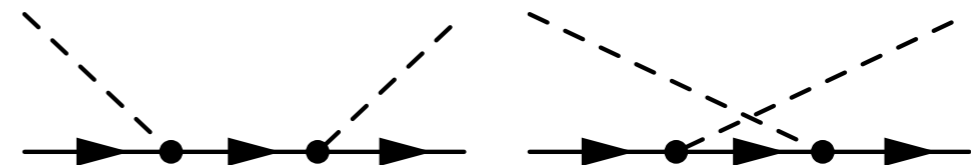
- Non-trivial result:
best NLO fit prefers **physical** values of **decay constants**:

f_K (MeV)	110.0	$(f_\pi = 92.4 \text{ MeV})$
f_η (MeV)	118.8	

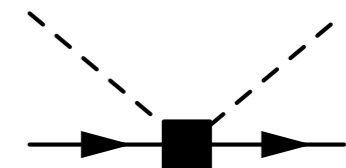
- **Tomozawa-Weinberg** terms **dominant**



- **Born** terms **significant**



- **NLO** parameters are non-negligible but **small**



UPDATED ANALYSIS of K^-p THRESHOLD PHYSICS with SIDDHARTA constraints (contd.)

	TW	TWB	NLO
$a_{\bar{K}N} (10^{-3})$	-1.57	-1.04	-2.38
$a_{\pi\Lambda} (10^{-3})$	-107.97	-8.06	-16.57
$a_{\pi\Sigma} (10^{-3})$	2.31	2.96	4.35
$a_{\eta\Lambda} (10^{-3})$	-0.20	-3.46	-0.01
$a_{\eta\Sigma} (10^{-3})$	216.37	3.52	1.90
$a_{K\Xi} (10^{-3})$	39.48	12.51	15.83
f_K (MeV)	110.8	109.0	110.0
f_η (MeV)	124.5	124.6	118.8
$\bar{b}_0 (10^{-2} \text{ GeV}^{-1})$	—	—	-4.79
$\bar{b}_D (10^{-2} \text{ GeV}^{-1})$	—	—	0.48
$\bar{b}_F (10^{-2} \text{ GeV}^{-1})$	—	—	4.01
$d_1 (10^{-2} \text{ GeV}^{-1})$	—	—	8.65
$d_2 (10^{-2} \text{ GeV}^{-1})$	—	—	-10.62
$d_3 (10^{-2} \text{ GeV}^{-1})$	—	—	9.22
$d_4 (10^{-2} \text{ GeV}^{-1})$	—	—	6.40
$\chi^2/\text{d.o.f.}$	1.12	1.15	0.96

Table 2: Parameters resulting from the systematic χ^2 analysis, using leading order (TW) plus Born terms (TWB) and full NLO schemes. Shown are the isospin symmetric subtraction constants $a_i(\mu)$ at $\mu = 1$ GeV, the meson decay constants f_K and f_η , the renormalized NLO constants \bar{b}_i and d_i , and $\chi^2/\text{d.o.f.}$ of the fit.

● Consistent LO \rightarrow NLO hierarchy

	TW	TWB	NLO
ΔE [eV]	373	377	306
Γ [eV]	495	514	591
γ	2.36	2.36	2.37
R_n	0.20	0.19	0.19
R_c	0.66	0.66	0.66
pole positions [MeV]	1422 - 16i 1384 - 90i	1421 - 17i 1385 - 105i	1424 - 26i 1381 - 81i

Table 3: Results of the systematic χ^2 analysis using leading order (TW) plus Born terms (TWB) and full NLO schemes. Shown are the energy shift and width of the 1s state of kaonic hydrogen (ΔE and Γ), threshold branching ratios (γ , R_n and R_c), and the pole positions of the isospin $I = 0$ amplitude in the $\bar{K}N-\pi\Sigma$ domain.

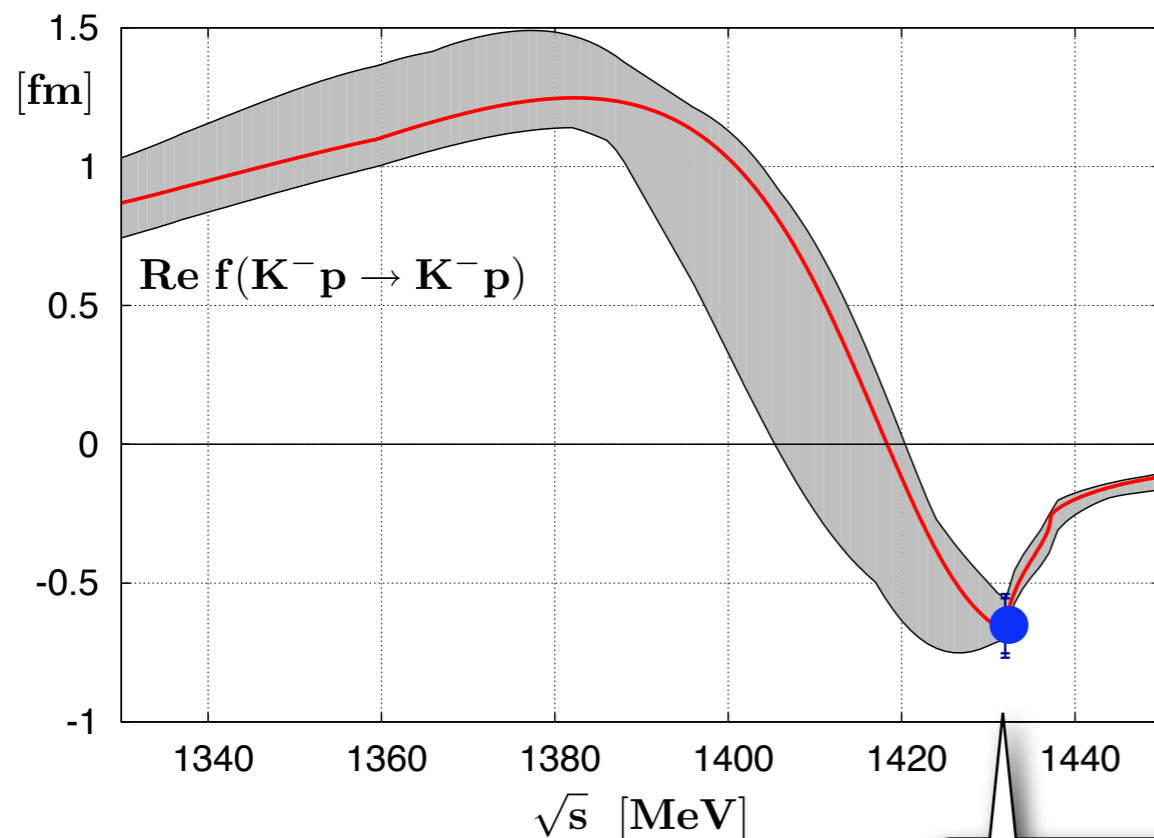
Y. Ikeda, T. Hyodo, W.W.
Physics Letters B 706 (2011) 63
Nucl. Phys. A 881 (2012) 98

$K^- p$ SCATTERING AMPLITUDE

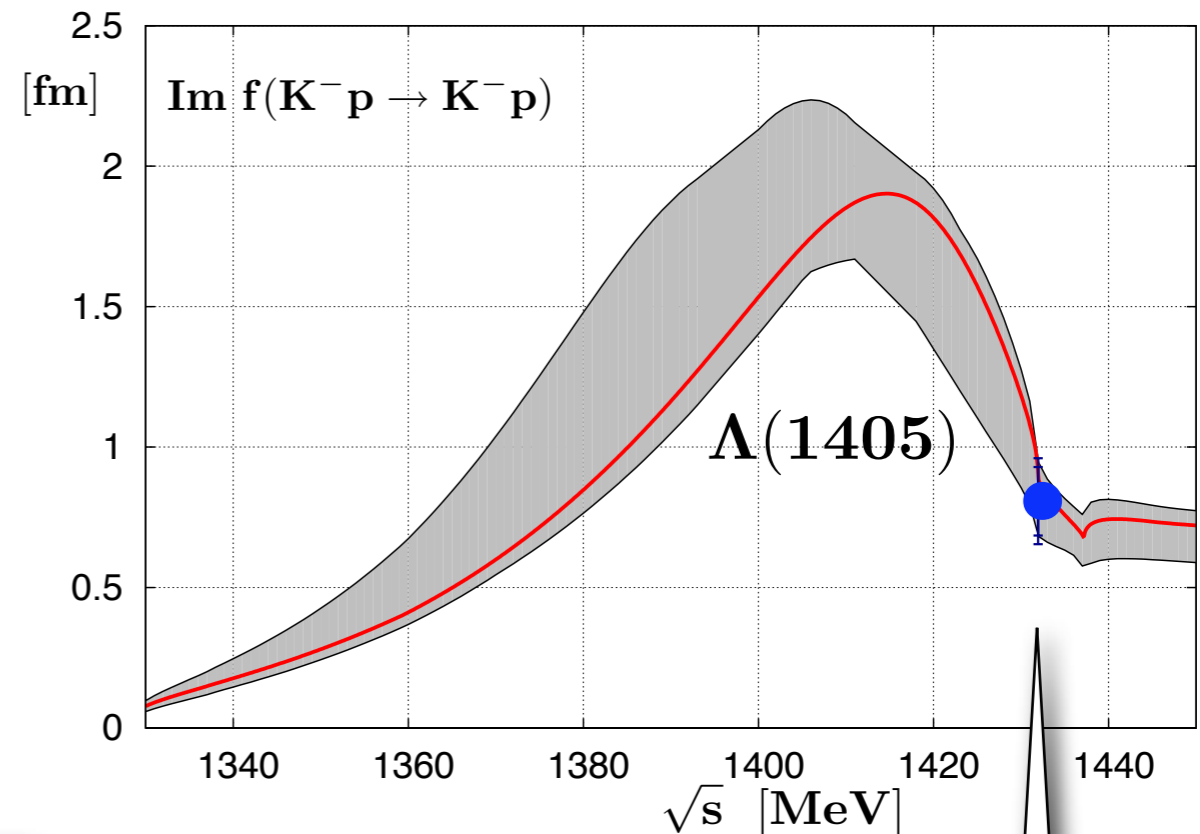
$$f(K^- p) = \frac{1}{2} [f_{\bar{K}N}(I=0) + f_{\bar{K}N}(I=1)]$$

- threshold region and subthreshold extrapolation:

$\Lambda(1405)$: $\bar{K}N$ ($I=0$) quasibound state embedded in the $\pi\Sigma$ continuum



$\text{Re } a(K^- p)$



$\text{Im } a(K^- p)$

- complex scattering length (including Coulomb corrections)

$$\text{Re } a(K^- p) = -0.65 \pm 0.10 \text{ fm}$$

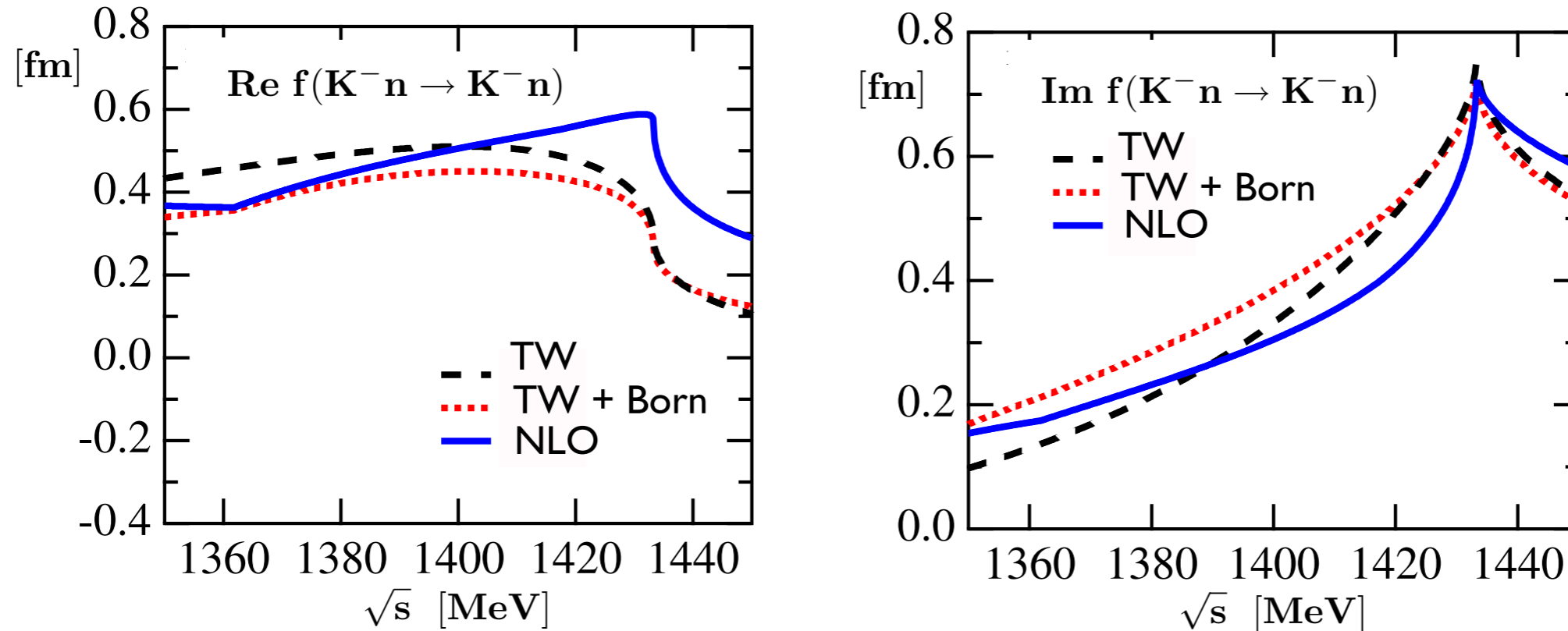
$$\text{Im } a(K^- p) = 0.81 \pm 0.15 \text{ fm}$$



CHIRAL SU(3) COUPLED CHANNELS DYNAMICS

- Predicted antikaon-neutron amplitudes at and below threshold

Y. Ikeda, T. Hyodo, W. Weise : Phys. Lett. B 706 (2011) 63 , Nucl. Phys. A 881 (2012) 98



$$a(K^- n) = 0.57^{+0.04}_{-0.21} + i 0.72^{+0.26}_{-0.41} \text{ fm}$$

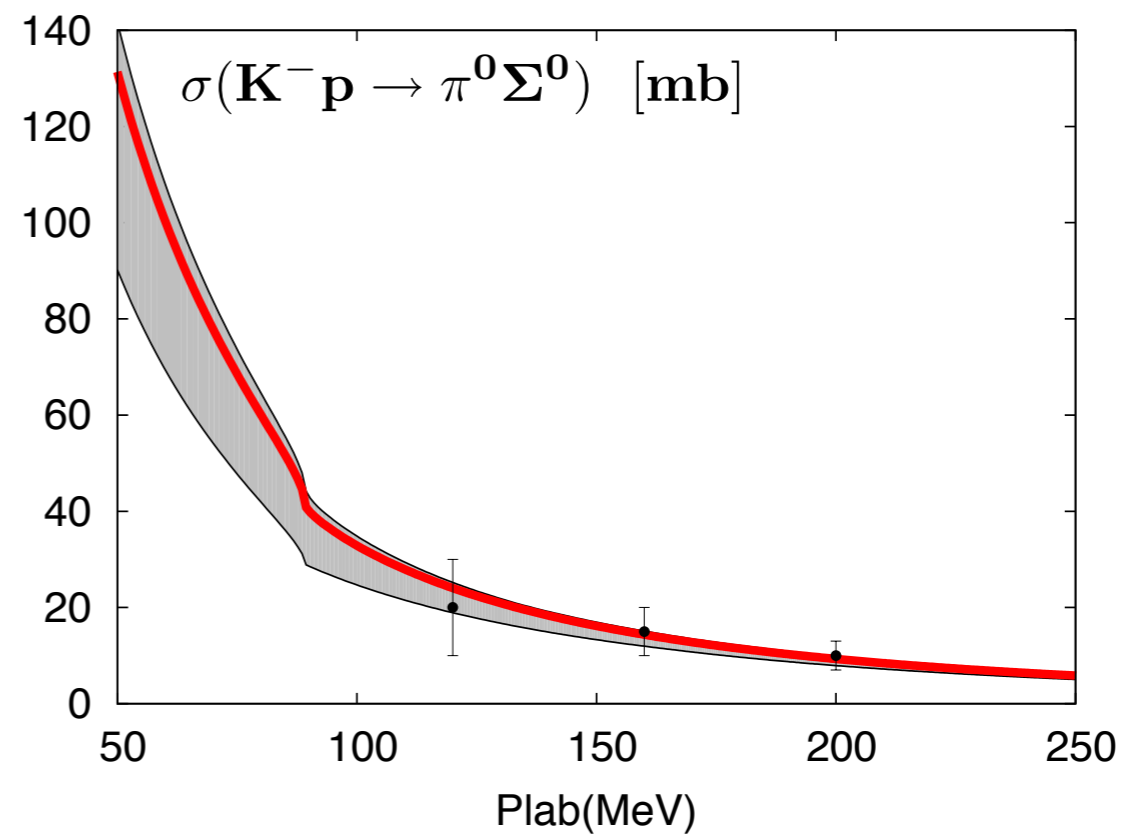
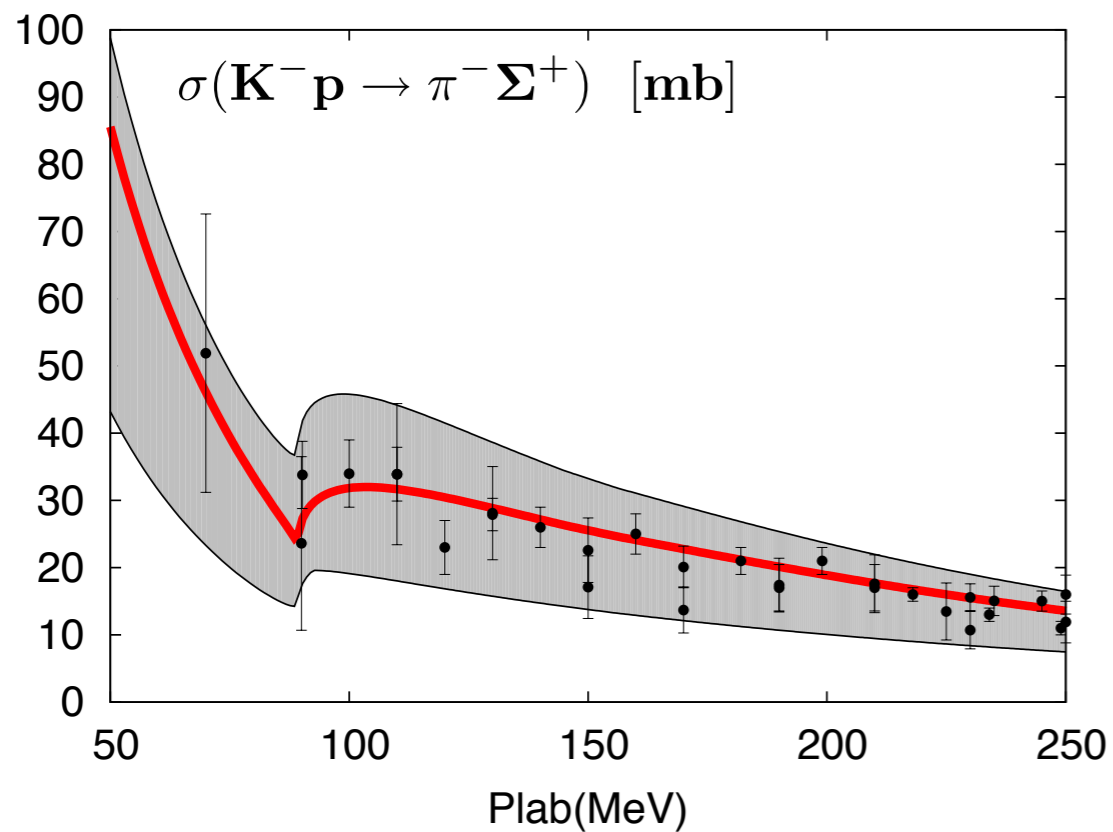
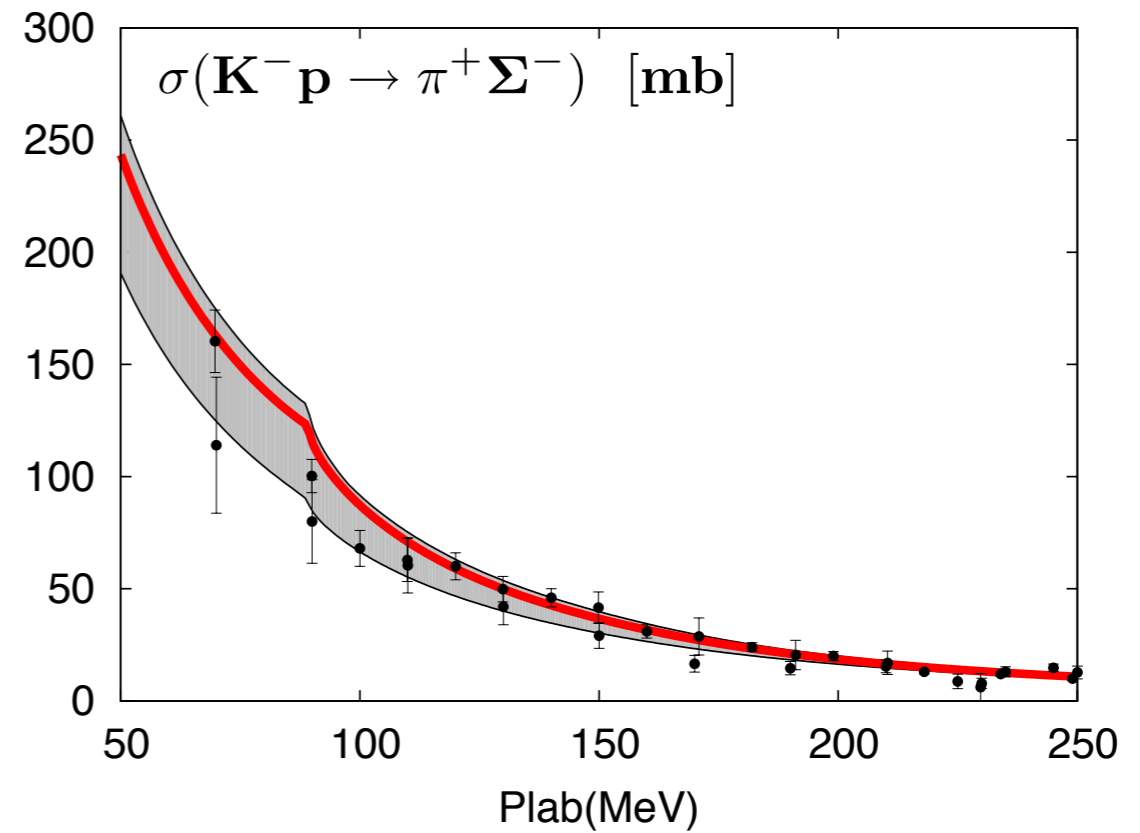
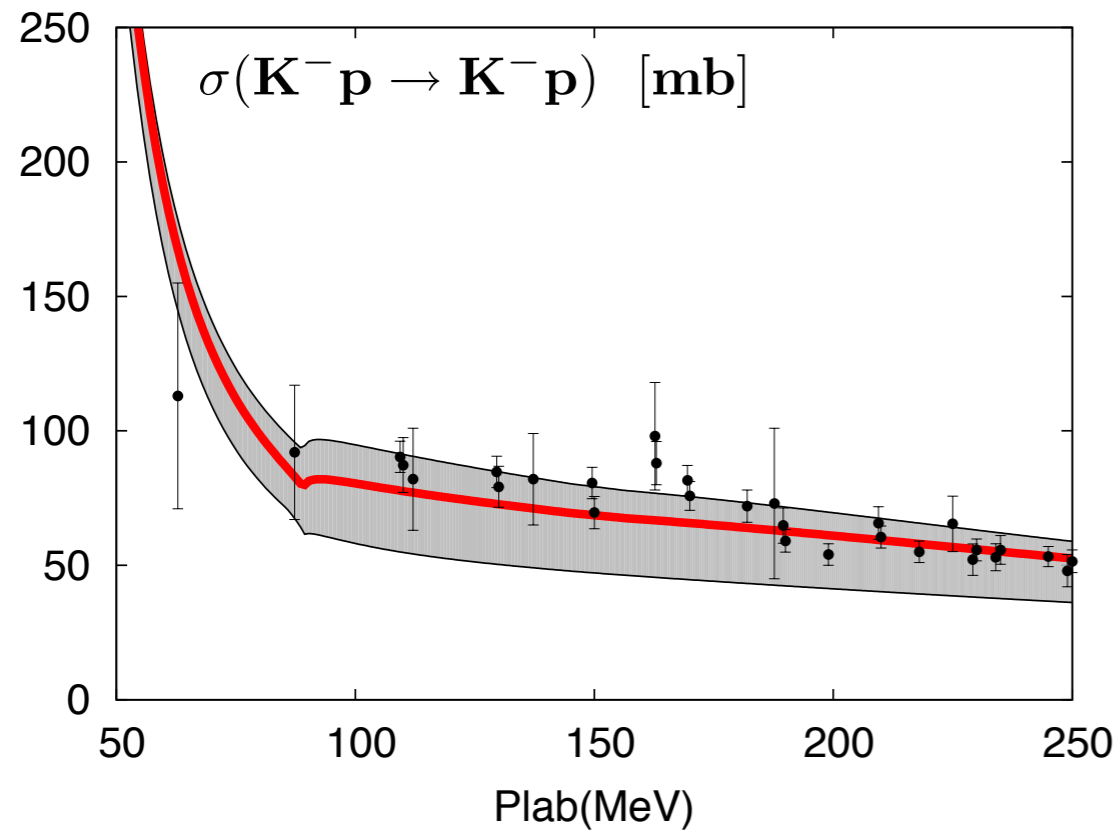
- **Needed:**

accurate constraints from antikaon-deuteron threshold measurements

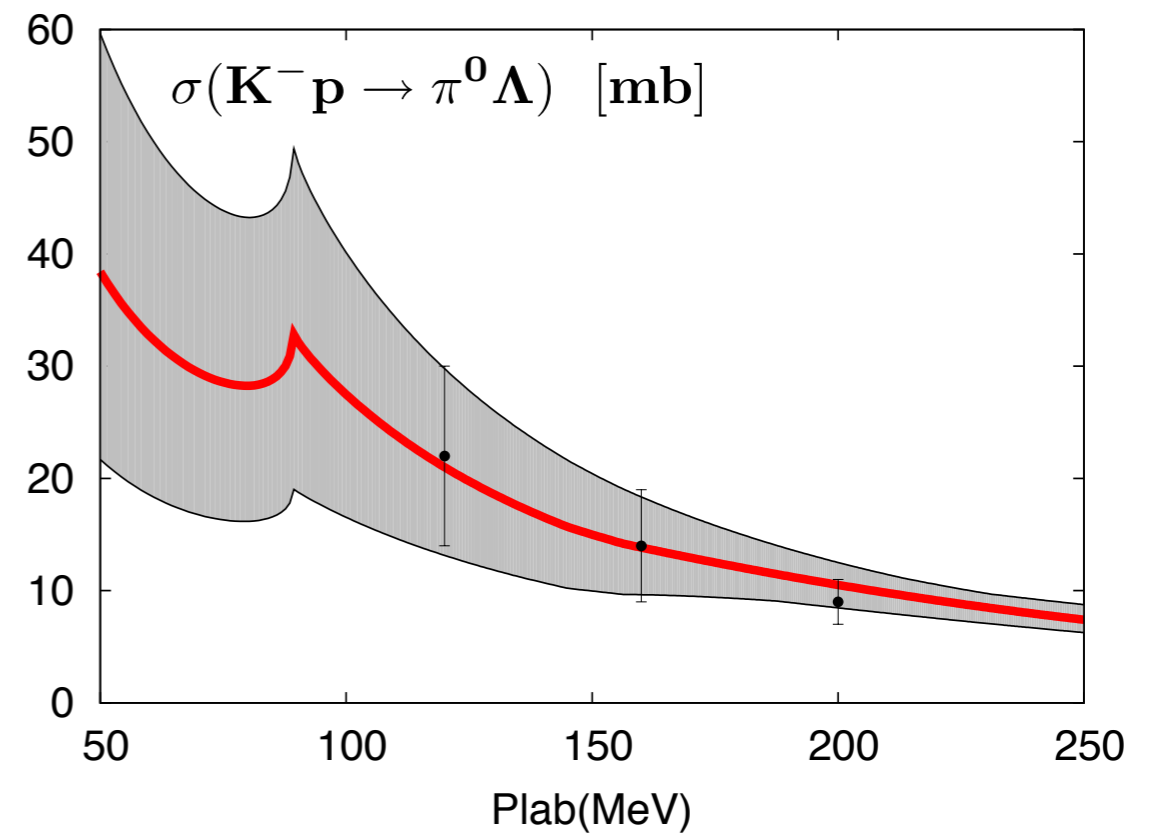
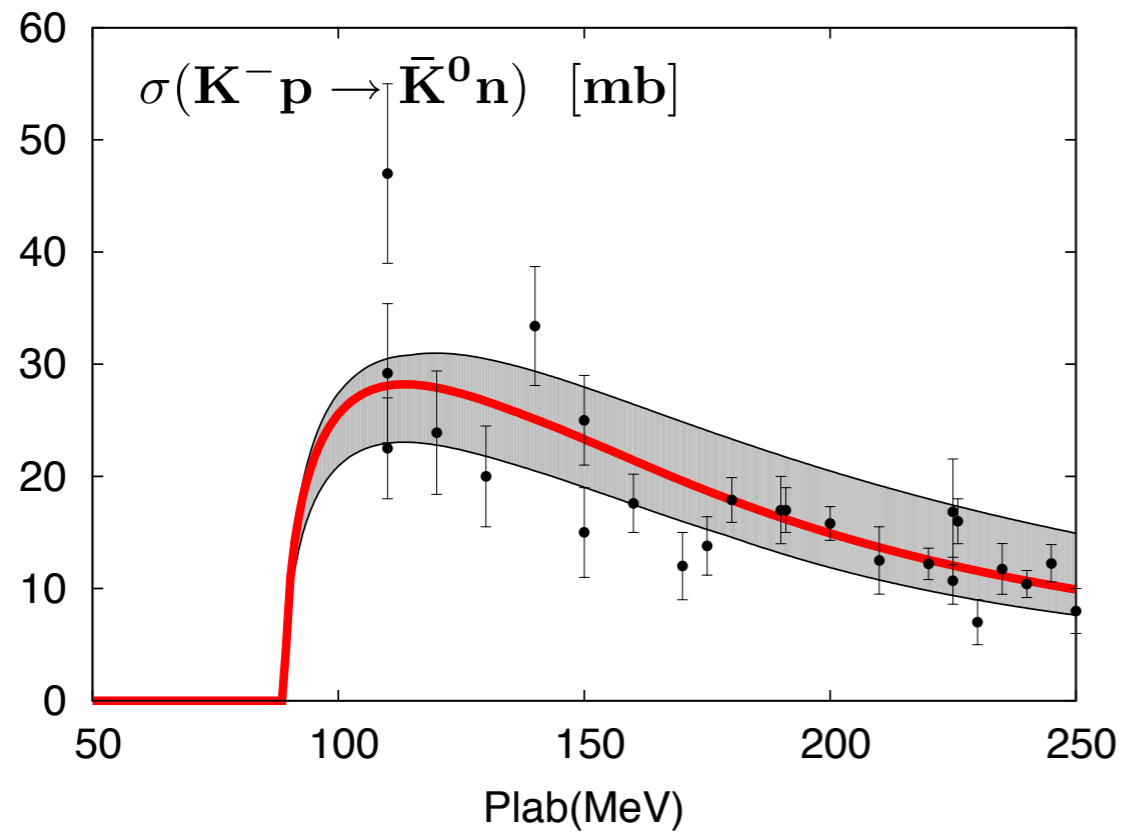
- ▶ **complete** information for both isospin $I = 0$ and $I = 1$ $\bar{K}N$ channels
- ▶ first measurement of **K-deuteron** scattering length
(including potentially important information about **K-NN absorption**)



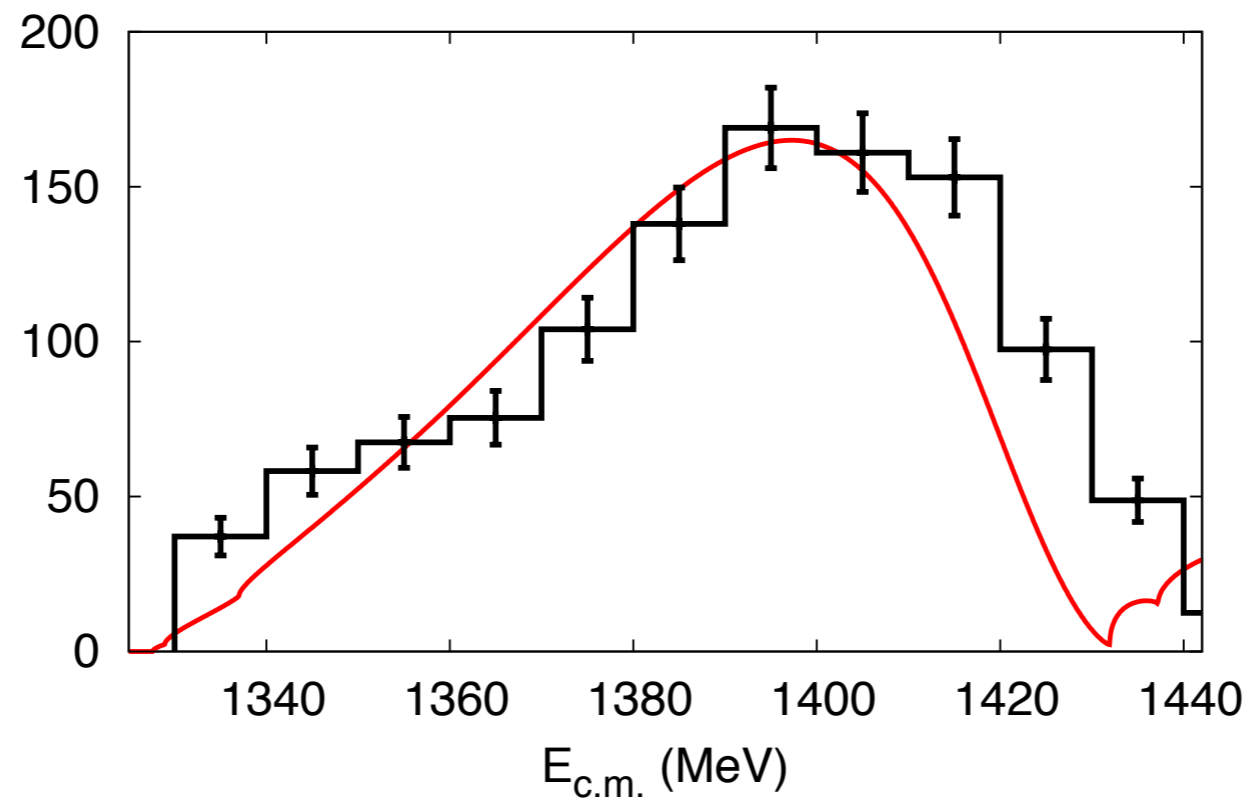
UPDATED ANALYSIS of K^-p LOW-ENERGY CROSS SECTIONS



UPDATED ANALYSIS of K^-p LOW-ENERGY CROSS SECTIONS



Mass spectrum $l=0 \pi\Sigma$

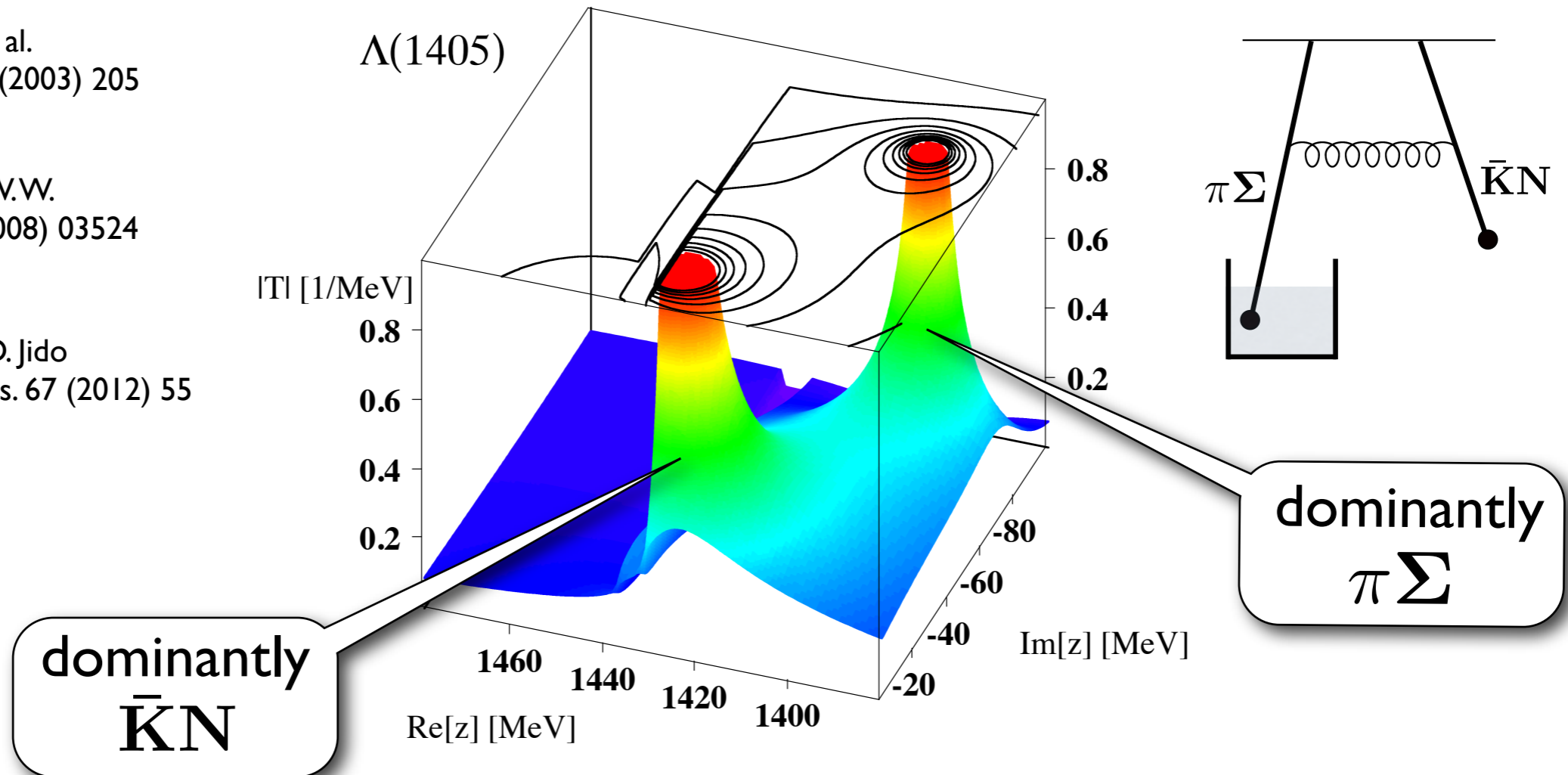


The TWO POLES scenario

D. Jido et al.
Nucl. Phys. A723 (2003) 205

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Phys. Rev. C 77 (2008) 03524

T. Hyodo, D. Jido
Prog. Part. Nucl. Phys. 67 (2012) 55



- Pole positions from chiral SU(3) coupled-channels calculation with SIDDHARTA threshold constraints:

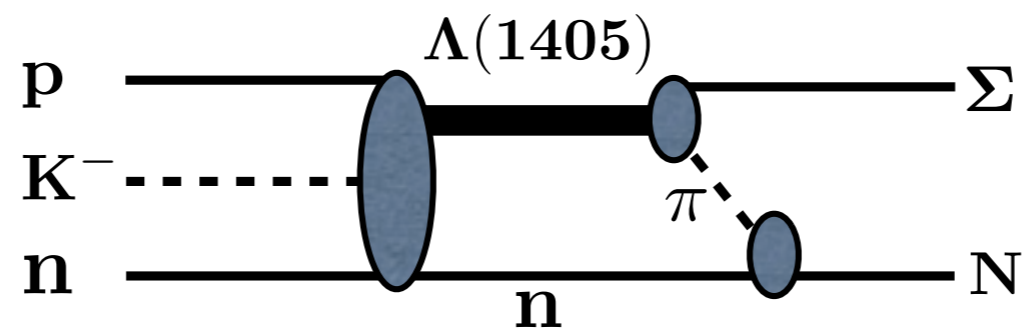
$E_1 = 1424 \pm 15 \text{ MeV}$	$E_2 = 1381 \pm 15 \text{ MeV}$
$\Gamma_1 = 52 \pm 10 \text{ MeV}$	$\Gamma_2 = 162 \pm 15 \text{ MeV}$

Y. Ikeda, T. Hyodo, W.W.: Nucl. Phys. A 881 (2012) 98

Implications & Comments

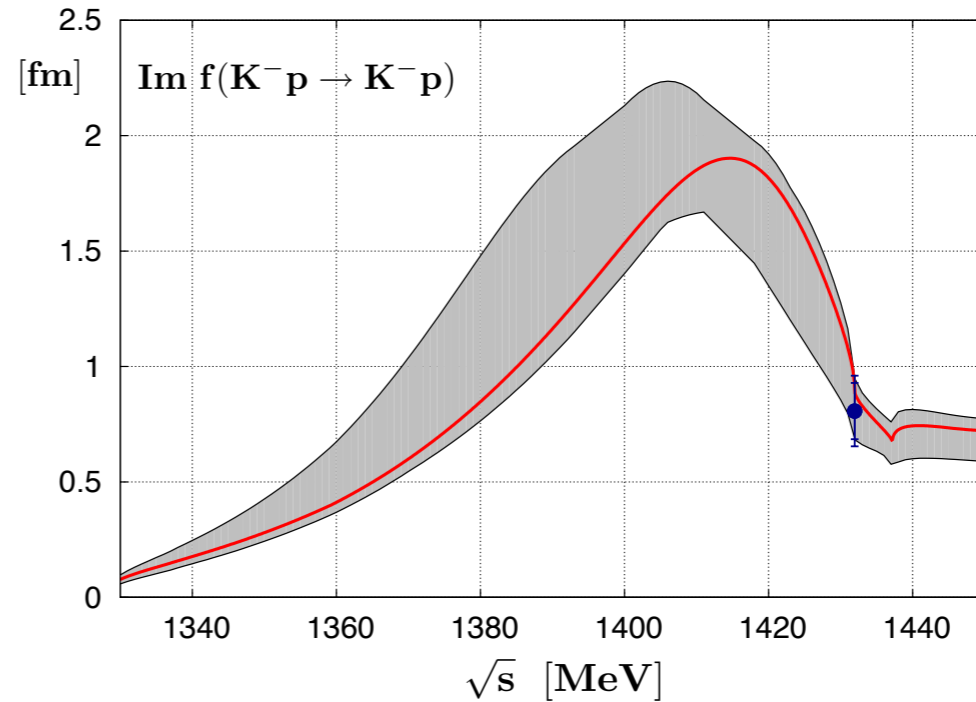
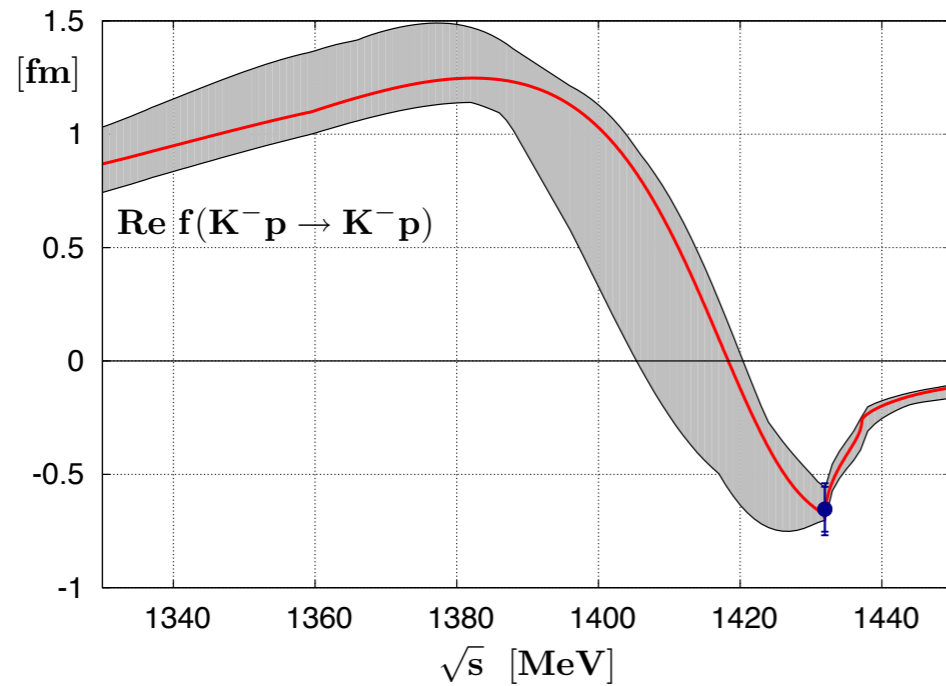
- $K^- p$ scattering length more accurately determined than $K^- n$ (SIDDHARTA constraints)
- Uncertainties in $\bar{K}N$ ($I = 1$) interaction primarily from large uncertainties in the $K^- p \rightarrow \pi^0 \Lambda$ channel
- **Kaonic deuterium** measurements important for providing further constraints on $K^- n$ interaction
- $B = 2$ systems - key issue: $\bar{K}NN \rightarrow YN$ absorption into non-mesonic hyperon-nucleon final states

e.g.:



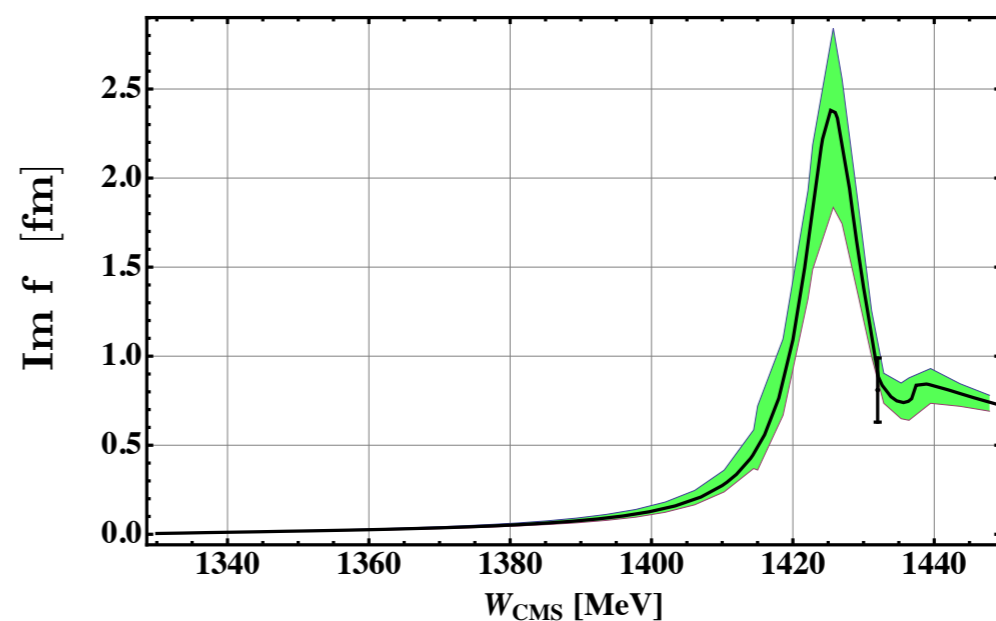
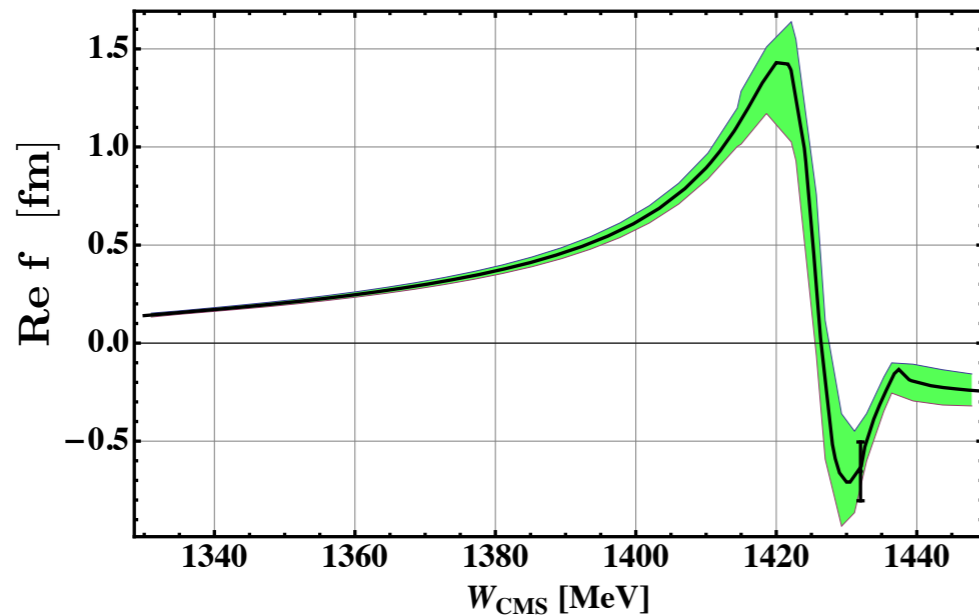
ALTERNATIVE OPTIONS ?

- Reproducing **kaonic hydrogen** and **low-energy scattering** data does not give **unique** answer - **subthreshold** constraints important



Y. Ikeda,
T. Hyodo,
W. W.

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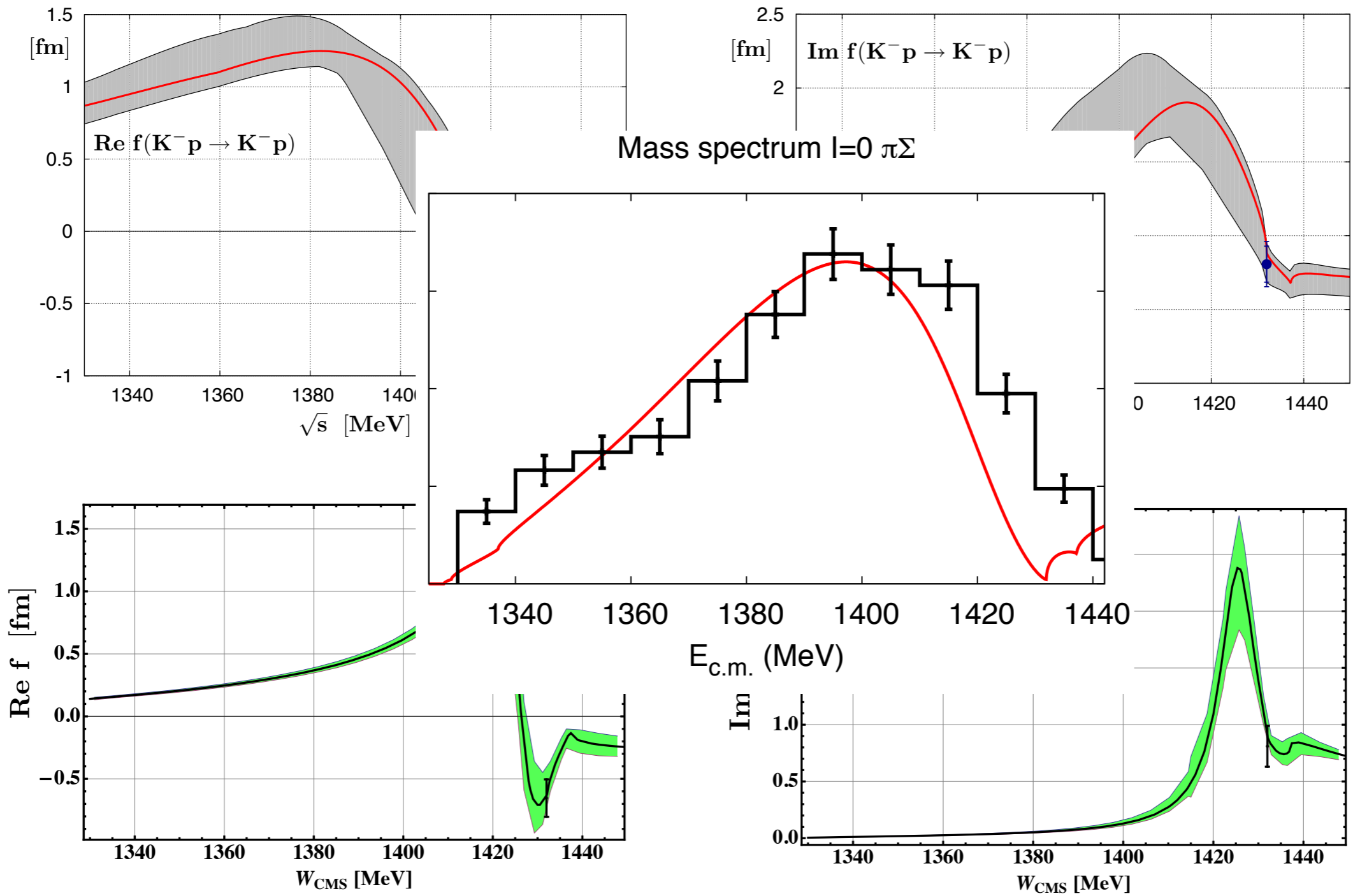
M. Mai,
U.-G. Meißner

arXiv:
1202.2030
[nucl-th]



ALTERNATIVE OPTIONS ?

- Reproducing **kaonic hydrogen** and **low-energy scattering** data does not give **unique** answer - **subthreshold** constraints important



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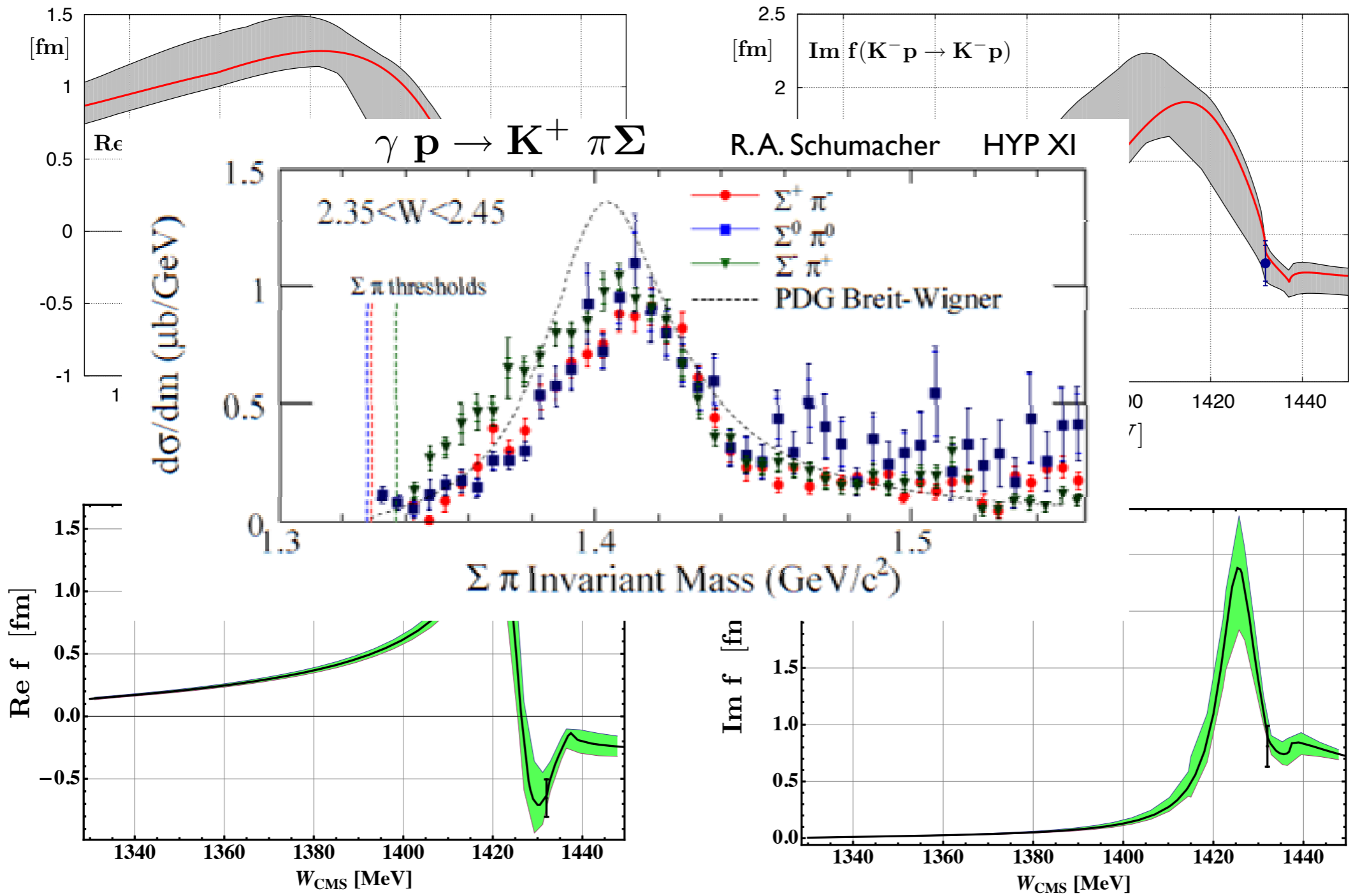
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1202.2030
[nucl-th]



ALTERNATIVE OPTIONS ?

- Reproducing **kaonic hydrogen** and **low-energy scattering** data does not give **unique** answer - **subthreshold** constraints important



Y. Ikeda,
T. Hyodo,
W. W.

PLB 706 (2011) 63
NPA881 (2012) 98

M. Mai,
U.-G. Meißner

arXiv:
1202.2030
[nucl-th]

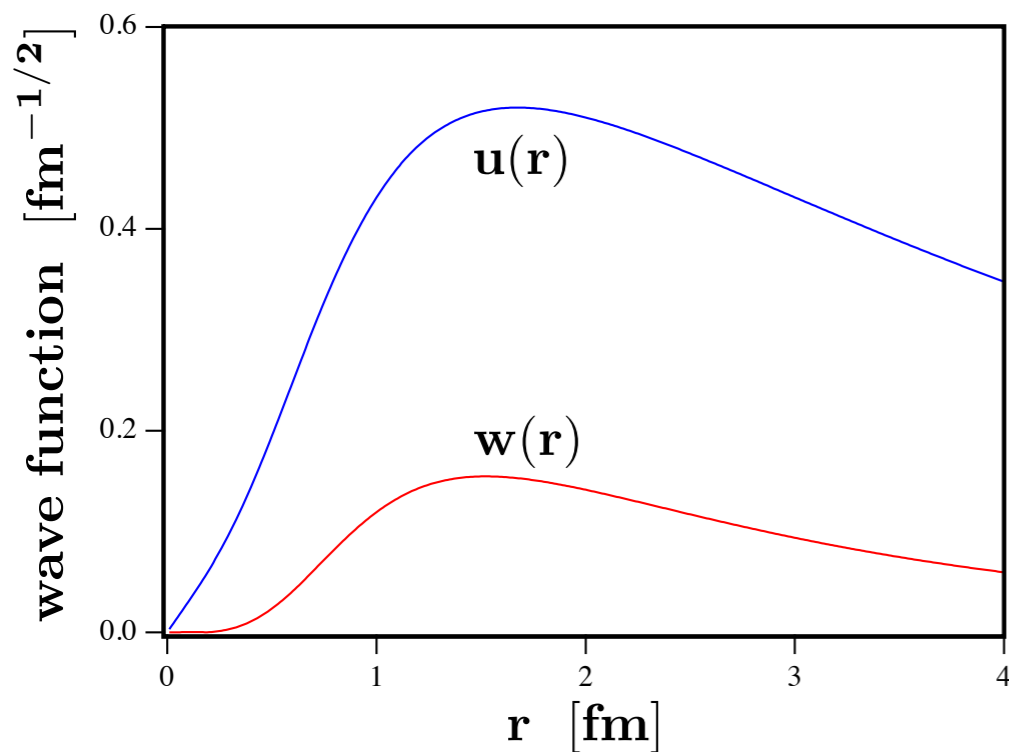


ANTIKAON - DEUTERON THRESHOLD PHYSICS

... looking forward to **SIDDHARTA 2**

- **Strategies:** Multiple scattering (MS) theory vs. three-body (Faddeev) calculations with Chiral SU(3) Coupled Channels input
- **MS approach** (fixed scatterer approximation): K^- d **scattering length**

S.S. Kamalov, E. Oset, A. Ramos: Nucl. Phys. A 690 (2001) 494



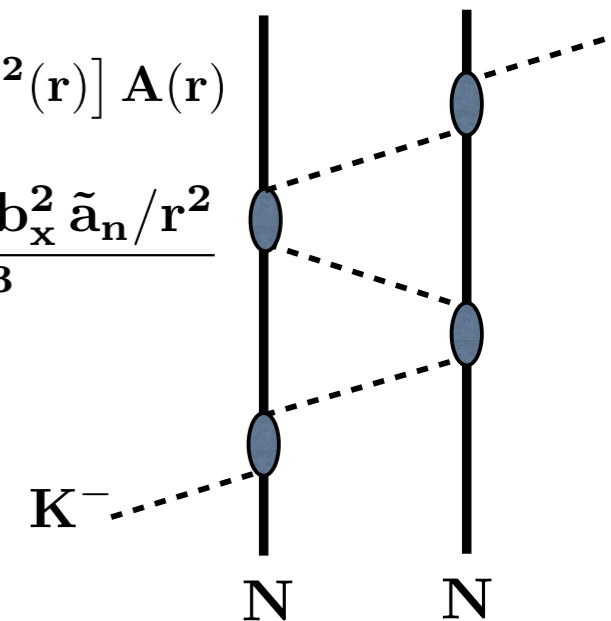
$$a(K^-d) = \left(1 + \frac{m_K}{M_d}\right)^{-1} \int_0^\infty dr [u^2(r) + w^2(r)] A(r)$$

$$A(r) = \frac{\tilde{a}_p + \tilde{a}_n + (2\tilde{a}_p \tilde{a}_n - b_x^2)/r - 2b_x^2 \tilde{a}_n/r^2}{1 - \tilde{a}_p \tilde{a}_n/r^2 + b_x^2 \tilde{a}_n/r^3}$$

$$\tilde{a}_p = \left(1 + \frac{m_K}{M_N}\right) a(K^-p \rightarrow K^-p)$$

$$\tilde{a}_n = \left(1 + \frac{m_K}{M_N}\right) a(K^-n \rightarrow K^-n)$$

b_x incorporates $K^-p \rightarrow \bar{K}_0 n$ and $\bar{K}_0 n \rightarrow \bar{K}_0 n$



- Using IHW input scattering lengths constrained by SIDDHARTA kaonic hydrogen:

$a(K^-d)$ [fm]

full MS	$-1.54 + i1.64$
no charge exchange	$-1.04 + i1.34$
impulse approximation	$-0.13 + i1.81$

T. Hyodo, Y. Ikeda, W.W.
(2012) preliminary

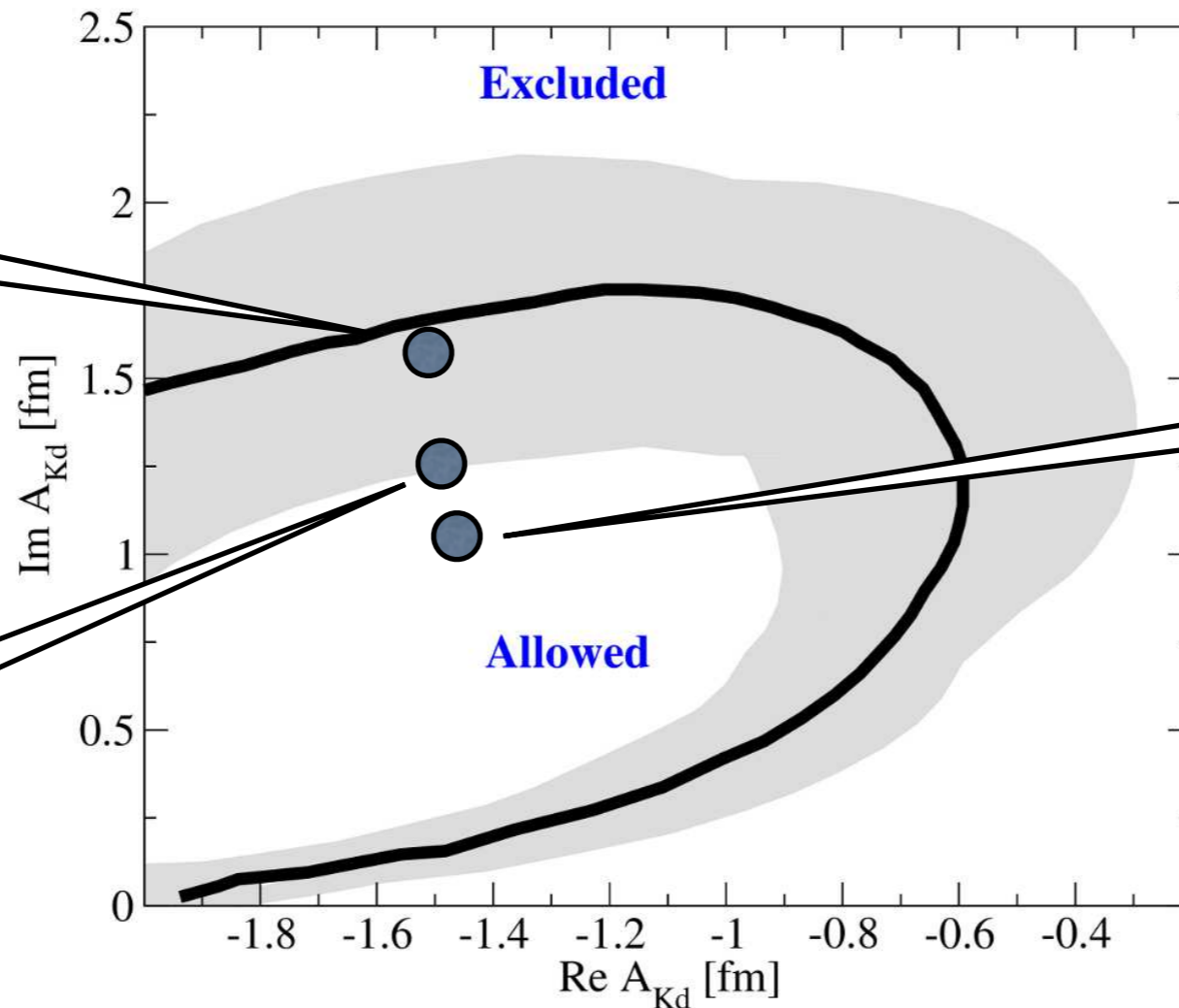


ANTIKAON - DEUTERON SCATTERING LENGTH

- Recent calculations using SIDDHARTA - constrained input

MS using IHW
scattering lengths

T. Hyodo, Y. Ikeda, W.W.
(2012) preliminary



Non-relativistic
effective field theory

M. Döring, U.-G. Meißner
Phys. Lett. B 704 (2011) 663

Faddeev calculation
separable “chiral”
amplitudes

N.V. Shevchenko
NPA 890-891 (2012) 50

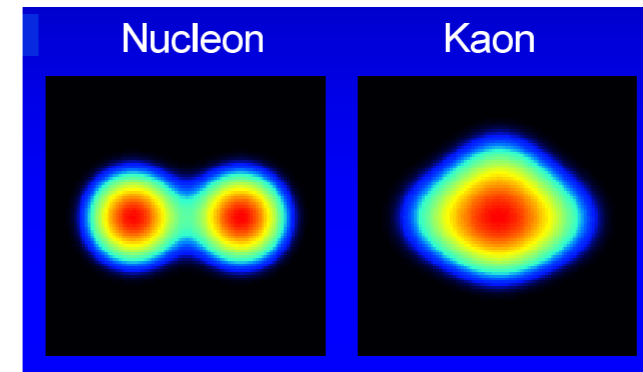
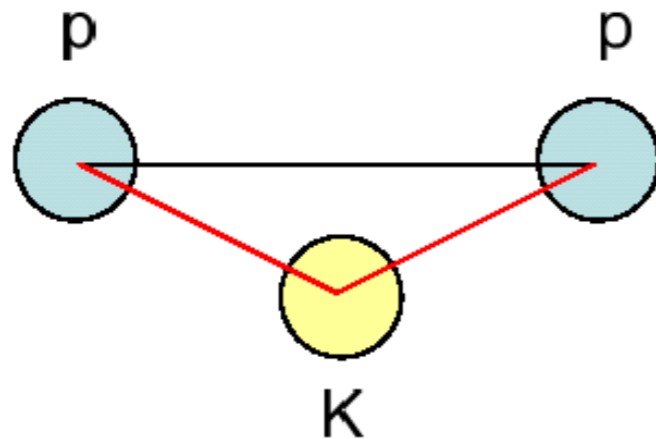
- Primary theoretical uncertainties from K^-n amplitude
- Predicted energy **shift** and **width** of kaonic deuterium (Faddeev calculation):

$$\Delta E_{1s} = -794 \text{ eV} \quad \Gamma_{1s}(K^-d) = 1012 \text{ eV}$$

- Not included: $K^-d \rightarrow YN$ absorption



UPDATE on QUASIBOUND K^-pp



3-Body (Faddeev) calculations

Variational calculations

- ... now consistently using amplitudes from **Chiral SU(3) coupled-channels** dynamics including **energy dependence** in subthreshold extrapolations

- Calculated **binding energy** and **width** (in MeV) of the K^-pp system

modest binding
large width

	[1]	[2]	[3]
B	16	17–23	9–16
Γ	41	40–70	34–46

remarkable degree of consistency

[1] Variational (hyperspherical harmonics): N. Barnea, A. Gal, E.Z. Livets ; Phys. Lett. B 712 (2012) 132

[2] Variational (Gaussian trial wave functions): A. Doté, T. Hyodo, W.W.; Phys. Rev. C 79 (2009) 014003

[3] Faddeev: Y. Ikeda, H. Kamano, T. Sato ; Prog. Theor. Phys. 124 (2010) 533

